

# Section 1: How QM Works, Part 1

In these slides we will cover:

- The Schrödinger Equation
- The probability interpretation of the wavefunction
- The discrete nature of observables
- The correspondence between observables and operators
- Eigenfunctions, eigenvalues and their properties
- Measurement in Quantum Mechanics
- Expectation values

# The wavefunction

## Particles and waves

- In classical physics, we use **Newton's Laws** to determine the equation of motion  $x(t)$  of a particle of mass  $m$  moving in a potential  $V(x)$ :

$$F = m \frac{d^2x}{dt^2} = - \frac{dV}{dx}$$

- Equivalently, we can **conserve the energy** of the particle:

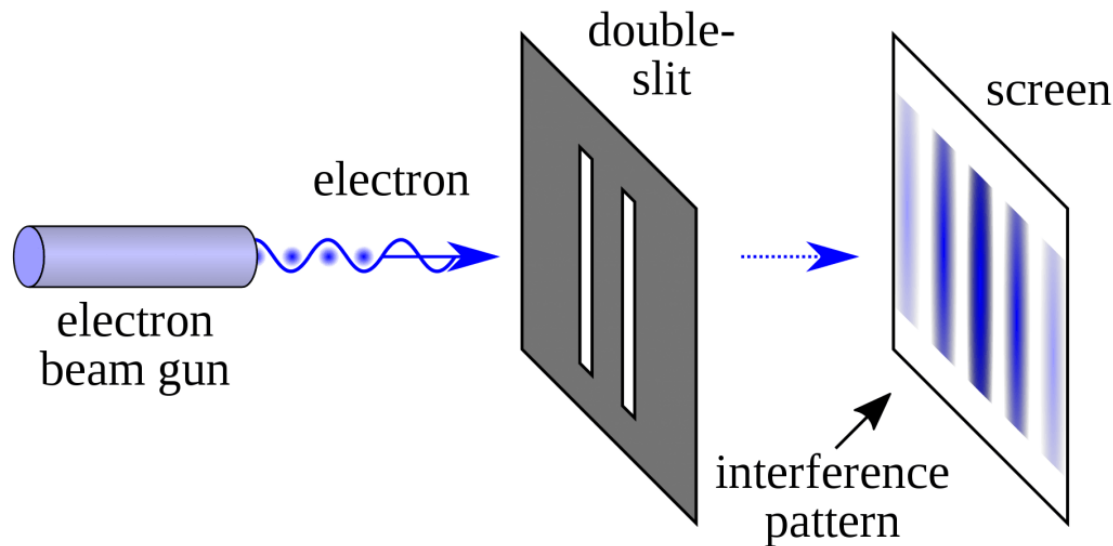
$$\frac{1}{2}mv^2 + V(x) = \text{Energy}$$



# The wavefunction

## Particles and waves

- This picture cannot apply in the Quantum world, because **particles behave like waves** (see: the double-slit experiment)



- Since a wave is an **object extended in space**, we need to change how we describe a particle

# The wavefunction

## The Schrödinger equation

- In Quantum Mechanics, the equation of motion of a particle in a potential  $V(x)$  is replaced by the **Schrödinger equation**:

$$-\frac{\hbar^2}{2m} \frac{\partial^2 \Psi(x, t)}{\partial x^2} + V(x) \Psi(x, t) = i\hbar \frac{\partial \Psi(x, t)}{\partial t}$$

- The symbol  $\hbar = h/2\pi$ , where  $h$  is Planck's constant
- It's an equation for the **wavefunction of the particle**  $\Psi(x, t)$ . This looks complicated, but we'll soon see it's the same as:

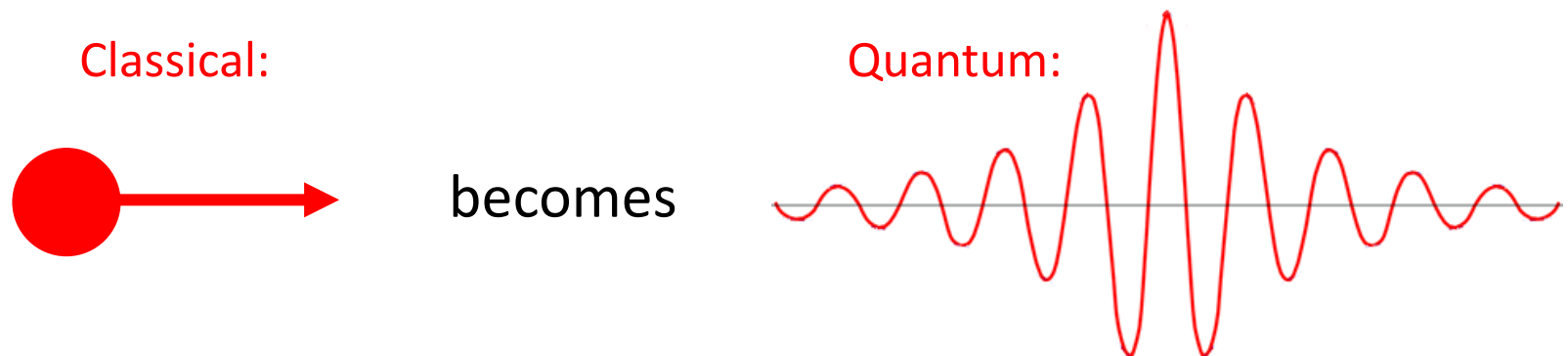
**Kinetic energy + Potential energy = Total energy**

- The  $i = \sqrt{-1}$  appearing in the Schrödinger equation looks strange – the wavefunction is a complex number in general!

# The wavefunction

## The wavefunction

- The **wavefunction**  $\Psi$  – that’s the Greek letter “psi” – is how we describe the state of a particle in Quantum Mechanics
- At a given time  $t$ , a particle is not at a fixed position  $x(t)$ , but is in a state described as a **function of position**,  $\Psi(x, t)$



- The wavefunction depends on the co-ordinates of a system and contains **all the information about the system**

# The wavefunction

## Probability interpretation of the wavefunction

- **What does the wavefunction mean?** It's connected to the **probability** of the particle being in a particular position:

$$\text{Probability of finding a particle in a range } x \rightarrow x + dx = |\Psi|^2 dx$$

Note: although  $\Psi$  can be a complex number,  $|\Psi|^2 = \Psi \cdot \Psi^*$  is real, as it should be for a probability!

- The particle must be somewhere! Hence, these probabilities must sum to 1.0, which is known as the **normalisation** of  $\Psi$ :

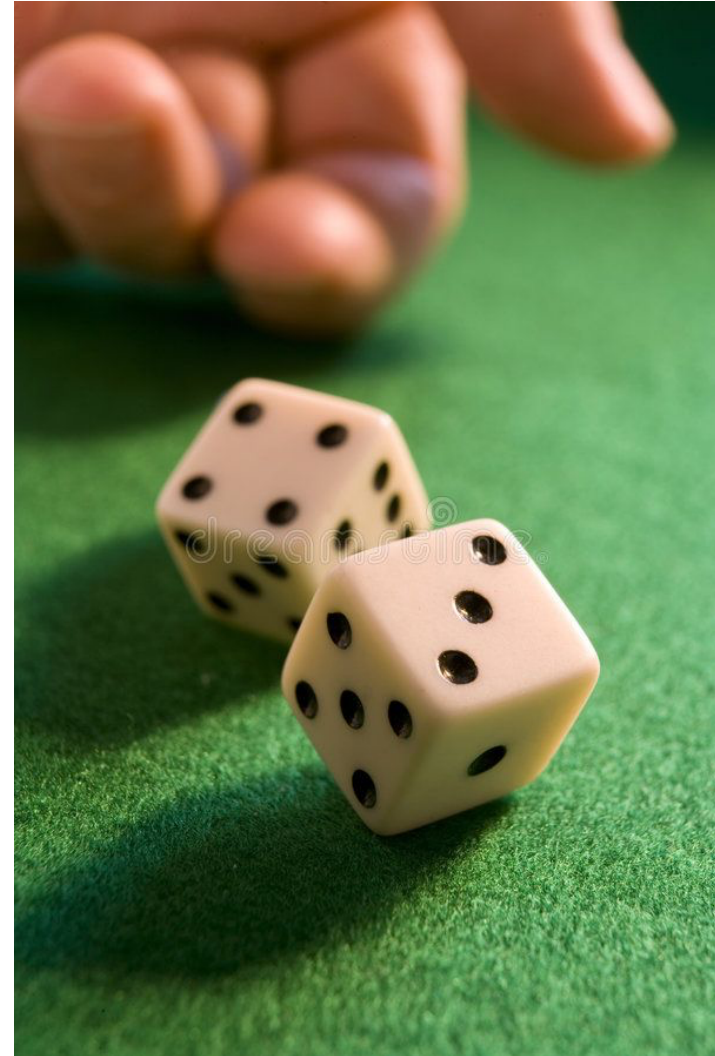
$$\int_{-\infty}^{\infty} |\Psi(x, t)|^2 dx = 1$$

- The probability interpretation of the wavefunction implies that Quantum Mechanics has a **statistical** or **indeterminate** nature

# Operators, eigenfunctions & eigenvalues

## Discrete nature of observables

- In Quantum Mechanics, a measurement of a quantity **can only produce discrete (specific) outcomes**, not any value
- You have previously studied a particle in an **infinite potential well**, which has certain allowed energy levels (see recap on next slide)
- Another example is poor Schrödinger's cat, which only has 2 possible states ...



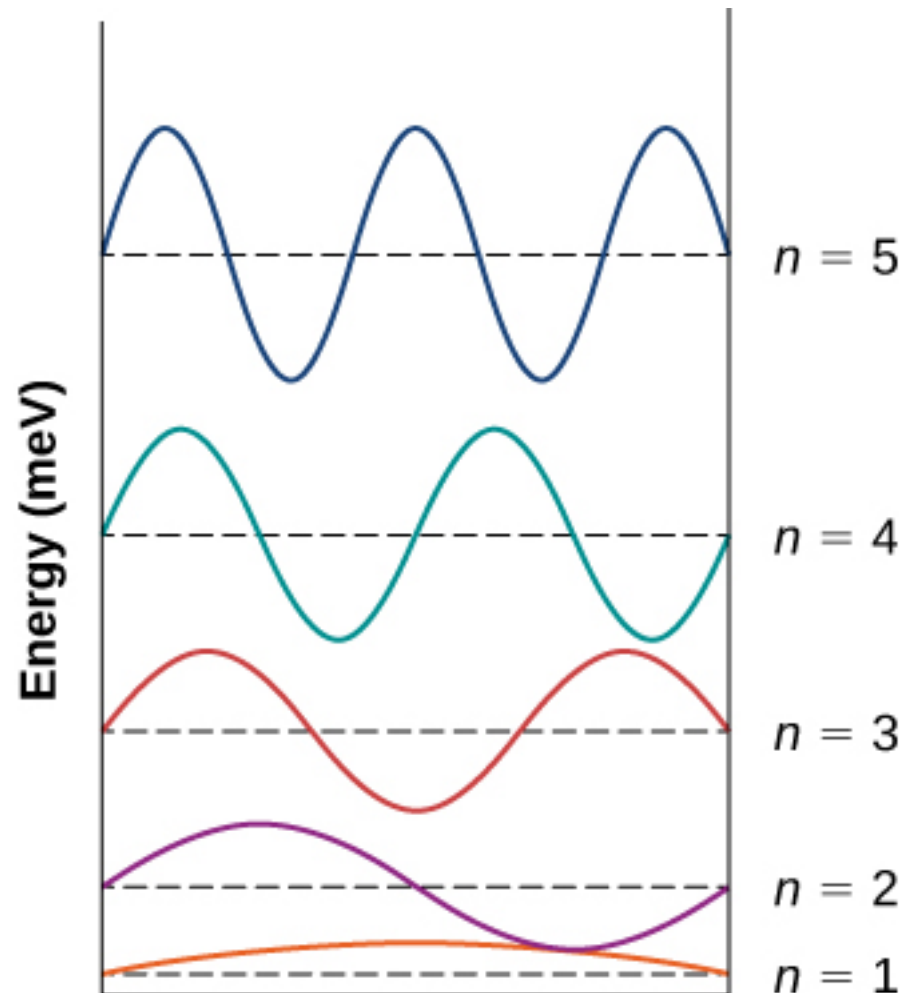
# Operators, eigenfunctions & eigenvalues

## Discrete nature of observables

- In Physics 2A QM, you studied that a particle enclosed in an infinite potential well has discrete energies and wavefunctions
- We'll see this example again in Section 3!

Image credit:

<https://opentextbc.ca/universityphysicsv3openstax/chapter/the-quantum-particle-in-a-box/>





# Operators, eigenfunctions & eigenvalues

## Discrete nature of observables

- Physics is described in the language of mathematics; so we need a mathematical structure in which discrete values appear
- Welcome to the world of **operators, eigenfunctions** and **eigenvalues**! Please do not turn back!
- We can describe the mathematical framework of Quantum Mechanics by the following statement:

*Each quantity we can observe is represented by a corresponding operator. If we measure that observable, we will always obtain a result which is one of the eigenvalues of the operator*

- What do these words mean??

# Operators, eigenfunctions & eigenvalues

## What is an operator?

- An operator is a **mathematical instruction which acts on a function to produce another function:**

**Operator** acts on **Function  $f(x)$**  to produce **Function  $g(x)$**

- **Example:**  $\frac{d}{dx}$  is an operator which acts on a function  $f(x)$  to produce the derivative function  $g(x) = \frac{df}{dx}$
- **Example:**  $x \cdot$  (“multiply by  $x$ ”) is an operator which acts on a function  $f(x)$  to produce another function  $g(x) = x f(x)$

# Operators, eigenfunctions & eigenvalues

## Eigenfunctions and eigenvalues

- When an operator acts on some special functions – called the **eigenfunctions** of the operator – it returns the same function, scaled by a number – called an **eigenvalue**

$$\hat{A} \phi_n(x) = a_n \phi_n(x)$$

$\hat{A}$  is an **operator** –  
these are usually  
written with little hats!

$\phi_n(x)$  is an  
**eigenfunction** – the  
subscript “ $n$ ” labels the  
different eigenfunctions  
( $\phi_1, \phi_2, \phi_3, \dots$ )

$a_n$  is the **eigenvalue**  
(number)  
corresponding to the  
eigenfunction  $\phi_n(x)$

# Operators, eigenfunctions & eigenvalues

## Eigenfunctions and eigenvalues

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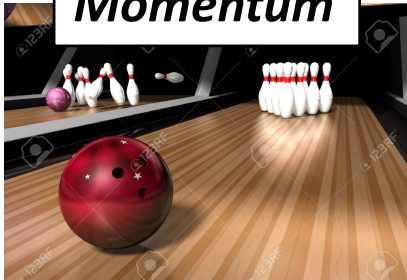
- As an example, let's consider the operator  $\hat{A} = \frac{d}{dx}$  again
- $\phi(x) = e^{ax}$  is an eigenfunction of  $\hat{A}$  with eigenvalue  $a$
- **Why?** Because  $\hat{A}\phi(x) = \frac{d\phi}{dx} = a e^{ax} = a \phi(x)$  – *the operator has returned the same function, scaled by a number*

# Operators, eigenfunctions & eigenvalues

## Properties of the operators representing observables

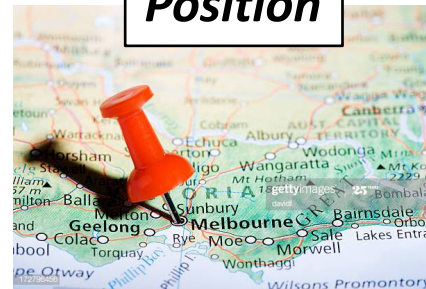
*Each quantity we can observe is represented by a corresponding operator. If we measure that observable, we will always obtain a result which is one of the eigenvalues of the operator*

### Momentum



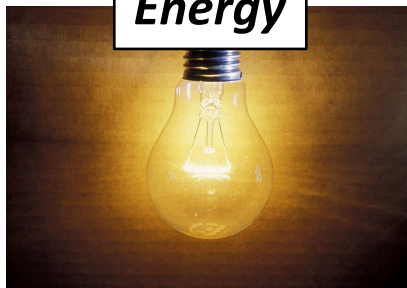
Represented by  
mathematical  
operator

### Position



Represented by  
mathematical  
operator

### Energy



Represented by  
mathematical  
operator

### Angular momentum



Represented by  
mathematical  
operator

# Operators, eigenfunctions & eigenvalues

## Properties of the operators representing observables

- The operators representing observables have **3 key properties**:
  1. Their **eigenvalues are real** (not complex) numbers, so they can correspond to the results of physical measurements

2. Different eigenfunctions are **orthogonal**, which is defined by:

$$\int_{-\infty}^{\infty} \phi_m^*(x) \phi_n(x) dx = \begin{cases} 1, & m = n \\ 0, & m \neq n \end{cases} \quad \text{Note: } \phi^* \text{ means the complex conjugate of } \phi$$

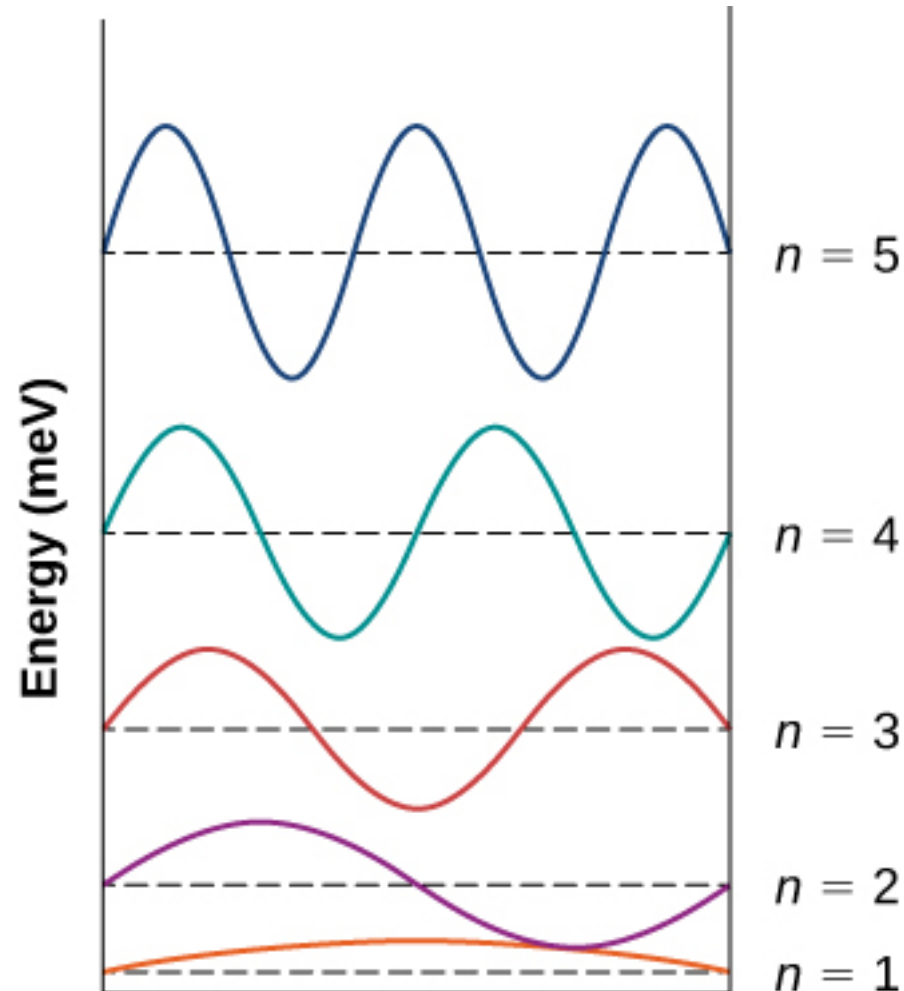
3. Any other function  $f(x)$  can be expressed as a **linear combination of the eigenfunctions**, which we can write as:

$$f(x) = \sum_n c_n \phi_n(x)$$

# Operators, eigenfunctions & eigenvalues

## Properties of the operators representing observables

- We can use the energy eigenfunctions for the infinite potential well to illustrate **orthogonality**
- These sine functions average to zero if  $m \neq n$ ,  
$$\int_{-\infty}^{\infty} \phi_m^*(x) \phi_n(x) dx = 0$$
- If  $m = n$ , then  
$$\int_{-\infty}^{\infty} |\phi_n(x)|^2 dx = 1$$
, which is the same as normalising the eigenfunctions



# Operators, eigenfunctions & eigenvalues

## Linear combinations of eigenfunctions

- We just mentioned that any function  $f(x)$  can be expressed as a **linear combination of the eigenfunctions** of an operator:

$$f(x) = \sum_n c_n \phi_n(x)$$

- We can determine the coefficients  $c_n$  using the orthogonality property. We can derive them by considering:

$$\int_{-\infty}^{\infty} \phi_m^*(x) f(x) dx = \int_{-\infty}^{\infty} \phi_m^*(x) \sum_n c_n \phi_n(x) dx$$

Changing the order of the integral and sum ...

$$= \sum_n c_n \int_{-\infty}^{\infty} \phi_m^*(x) \phi_n(x) dx$$

$$= c_m$$

This is equal to 1 if  $m = n$   
and 0 otherwise



# Measurement in Quantum Mechanics

## Measurement if $\Psi$ is an eigenfunction

- At the heart of Quantum Mechanics is the how the wavefunction is related to **measurement of observables**
- Suppose we measure a particular observable of a system (e.g. momentum, position, energy, angular momentum, etc.)
- Recapping ... this observable is represented by a corresponding **operator** (we will see some examples shortly)
- If the wavefunction of the system is an **eigenfunction** of the corresponding operator ( $\Psi = \phi_n$ ), then the result of the measurement is the corresponding **eigenvalue**,  $a_n$

# Measurement in Quantum Mechanics

## Measurement if $\Psi$ is not an eigenfunction

- If the wavefunction  $\Psi$  is **not** an eigenfunction of the corresponding operator, it can always be expressed as a **linear combination** of the eigenfunctions:

$$\Psi(x) = \sum_n c_n \phi_n(x)$$

- In this case, **the result of the measurement may be any one of the eigenvalues  $a_n$ , with corresponding probabilities  $|c_n|^2$**
- Following the measurement, **the wavefunction “collapses”** and becomes the eigenfunction,  $\Psi(x) = \phi_n(x)$
- If the observable is measured again, we'll find value  $a_n$  again

# Measurement in Quantum Mechanics

**Measurement if  $\Psi$  is not an eigenfunction**

*Measurement changes the wavefunction, causing it to “collapse” into the eigenfunction corresponding to the result of the measurement. This ensures that future measurements of the quantity produce the same result.*

- It is equivalent to opening the box containing Schrödinger’s cat, and seeing the result!

[Yes, cat photos are an occupational hazard here.]

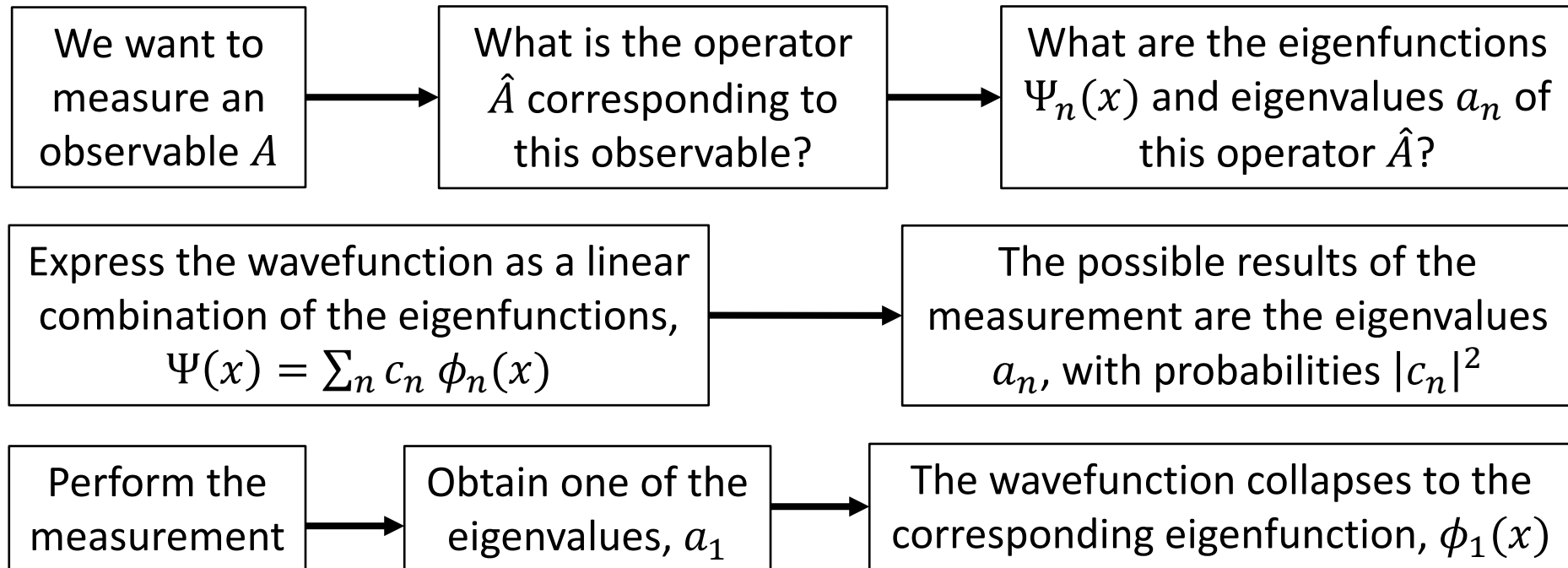


# Measurement in Quantum Mechanics

## Measurement if $\Psi$ is not an eigenfunction

- This gives us the **recipe for measurement in Quantum Mechanics**, represented by the following flow chart!

The state of the particle is described by its wavefunction  $\Psi(x)$



# Measurement in Quantum Mechanics

## Measurement if $\Psi$ is not an eigenfunction

*Example: a particle in an infinite potential well in the range  $|x| < L$  is prepared in the wavefunction  $\Psi(x) = \frac{1}{\sqrt{5L}} \left[ \cos\left(\frac{\pi x}{2L}\right) + 2 \sin\left(\frac{\pi x}{L}\right) \right]$ . What energy states can be measured, and with what probabilities?*

- The factor  $\frac{1}{\sqrt{5L}}$  is to ensure  $\Psi(x)$  is normalised:  $\int_{-\infty}^{\infty} |\Psi(x)|^2 dx = 1$
- We notice that the terms in the square bracket are the 1<sup>st</sup> and 2<sup>nd</sup> energy eigenfunctions. Substituting these in, we find  $\Psi(x) = \frac{1}{\sqrt{5}} \phi_1(x) + \frac{2}{\sqrt{5}} \phi_2(x)$
- Hence, the possible measurements of the energy state are  $E_1$  (with probability  $|c_1|^2 = \left| \frac{1}{\sqrt{5}} \right|^2 = \frac{1}{5}$ ) and  $E_2$  (with probability  $|c_2|^2 = \left| \frac{2}{\sqrt{5}} \right|^2 = \frac{4}{5}$ )
- The probabilities  $\frac{1}{5}$  and  $\frac{4}{5}$  sum to 1, as they should!

# Measurement in Quantum Mechanics

## Expectation values

- Although we can't predict exactly which value will result from a measurement (only their probabilities), we can predict the **mean measurement**, also known as the **expectation value** – which has the symbol of angled brackets,  $\langle a \rangle$

$$\langle a \rangle = \sum_n \text{Prob}(n) a_n$$

- **Example:** what is the expectation value when a dice is thrown?

$$\langle a \rangle = \frac{1}{6} \cdot 1 + \frac{1}{6} \cdot 2 + \frac{1}{6} \cdot 3 + \frac{1}{6} \cdot 4 + \frac{1}{6} \cdot 5 + \frac{1}{6} \cdot 6 = 3.5$$

*[Note: the value 3.5 cannot be obtained for any individual throw of the dice, but is the average over many throws!]*



# Measurement in Quantum Mechanics

## Expectation values

- In Quantum Mechanics, we have seen that  $\text{Prob}(n) = |c_n|^2$

$$\langle a \rangle = \sum_n \text{Prob}(n) a_n = \sum_n |c_n|^2 a_n$$

- Example:** for the particle in the infinite potential well 2 slides back,  $\langle E \rangle = \frac{1}{5} E_1 + \frac{4}{5} E_2$
- We can also find a general relation by substituting in  $c_n = \int_{-\infty}^{\infty} \Psi(x) \phi_n^*(x) dx$  and using the orthogonality relation:

$$\langle a \rangle = \int_{-\infty}^{\infty} \Psi^*(x) \hat{A} \Psi(x) dx$$

# Summary

## The rules of Quantum Mechanics

- The state of a particle is described by the wavefunction  $\Psi(x, t)$ , which satisfies the Schrödinger equation, 
$$-\frac{\hbar^2}{2m} \frac{\partial^2 \Psi}{\partial x^2} + V(x) \Psi = i\hbar \frac{\partial \Psi}{\partial t}$$
- The probability of a particle being located at a position in the range  $[x, x + dx]$  at time  $t$  is  $|\Psi(x, t)|^2$
- Physical observables are represented by operators, and the possible results of measuring an observable are the eigenvalues  $a_n$  of those operators
- Any wavefunction can be expressed as a linear combination of the eigenfunctions of these operators:  $\Psi = \sum_n c_n \phi_n$ . The probability of measuring the eigenvalue  $a_n$  is then  $|c_n|^2$
- If a measurement of an observable yields a result  $a_n$ , the wavefunction collapses into the corresponding eigenfunction,  $\Psi = \phi_n$