

## Quantum Mechanics Week 4: Class Activities solutions

Q1)

a) The equation with  $l = 0$  and  $m = 0$  is:

$$\frac{1}{\sin \theta} \frac{d}{d\theta} \left( \sin \theta \frac{dP_{00}}{d\theta} \right) = 0$$

which is satisfied if  $P_{00} = 1$ , since  $\frac{dP_{00}}{d\theta} = 0$ .

b) The equation with  $l = 1$  and  $m = 0$  is:

$$\frac{1}{\sin \theta} \frac{d}{d\theta} \left( \sin \theta \frac{dP_{10}}{d\theta} \right) + 2P_{10} = 0$$

Substituting in  $P_{10} = \cos \theta$ , we find:

$$\begin{aligned} \frac{1}{\sin \theta} \frac{d}{d\theta} \left( \sin \theta \frac{d}{d\theta} \cos \theta \right) + 2 \cos \theta &= \frac{1}{\sin \theta} \frac{d}{d\theta} (-\sin^2 \theta) + 2 \cos \theta \\ &= \frac{1}{\sin \theta} (-2 \sin \theta \cos \theta) + 2 \cos \theta \\ &= 0 \end{aligned}$$

demonstrating that this form is a solution.

c) The equation with  $l = 1$  and  $m = 1$  is:

$$\frac{1}{\sin \theta} \frac{d}{d\theta} \left( \sin \theta \frac{dP_{11}}{d\theta} \right) + \left[ 2 - \frac{1}{\sin^2 \theta} \right] P_{11} = 0$$

Substituting in  $P_{11} = -\sin \theta$ , we find:

$$\begin{aligned} \frac{1}{\sin \theta} \frac{d}{d\theta} \left( \sin \theta \frac{d}{d\theta} (-\sin \theta) \right) + \left[ 2 - \frac{1}{\sin^2 \theta} \right] (-\sin \theta) \\ &= -\frac{1}{\sin \theta} \frac{d}{d\theta} (\sin \theta \cos \theta) - 2 \sin \theta + \frac{1}{\sin \theta} \\ &= -\frac{1}{\sin \theta} (\cos^2 \theta - \sin^2 \theta) - 2 \sin \theta + \frac{1}{\sin \theta} \\ &= -\frac{1}{\sin \theta} (1 - 2 \sin^2 \theta) - 2 \sin \theta + \frac{1}{\sin \theta} \\ &= 0 \end{aligned}$$

demonstrating that this form is a solution.

Hence, using the form  $Y_{lm}(\theta, \phi) = N P_{lm}(\theta) e^{im\phi}$ , the mathematical functions for  $Y_{00}$ ,  $Y_{10}$  and  $Y_{11}$  are (ignoring the normalisation constant for now):

$$\begin{aligned} Y_{00} &\propto 1 \\ Y_{10} &\propto \cos \theta \\ Y_{11} &\propto -\sin \theta e^{i\phi} \end{aligned}$$

Q2)

The “ $\sin \theta$ ” appears in the equation for normalising the spherical harmonic functions because it is part of the area element for integrations over a sphere. There is more area available near the equator than at the poles.

For  $Y_{00}$ :

$$\begin{aligned}\int_0^{2\pi} d\phi \int_0^\pi d\theta \sin \theta N^2 &= 1 \\ \rightarrow 2\pi \cdot [\cos \theta]_\pi^0 \cdot N^2 &= 1 \\ \rightarrow 2\pi \cdot 2 \cdot N^2 &= 1 \\ \rightarrow N &= \frac{1}{\sqrt{4\pi}}\end{aligned}$$

For  $Y_{10}$ :

$$\begin{aligned}\int_0^{2\pi} d\phi \int_0^\pi d\theta \sin \theta N^2 \cos^2 \theta &= 1 \\ \rightarrow 2\pi \cdot \left[ \frac{\cos^3 \theta}{3} \right]_\pi^0 \cdot N^2 &= 1 \\ \rightarrow 2\pi \cdot \frac{2}{3} \cdot N^2 &= 1 \\ \rightarrow N &= \sqrt{\frac{3}{4\pi}}\end{aligned}$$

For  $Y_{11}$  (note the  $e^{i\phi}$  cancels out in the complex conjugate):

$$\int_0^{2\pi} d\phi \int_0^\pi d\theta \sin \theta N^2 \sin^2 \theta = 1$$

We can solve the integral of  $\sin^3 \theta$  over  $\theta$  using the trig formula  $\sin 3\theta = 3 \sin \theta - 4 \sin^3 \theta$ , hence  $\sin^3 \theta = \frac{3}{4} \sin \theta - \frac{1}{4} \sin 3\theta$ . Integrating,

$$\begin{aligned}\rightarrow 2\pi \cdot \left[ \frac{3}{4} \cos \theta - \frac{1}{12} \cos 3\theta \right]_\pi^0 \cdot N^2 &= 1 \\ \rightarrow 2\pi \cdot \frac{4}{3} \cdot N^2 &= 1 \\ \rightarrow N &= \sqrt{\frac{3}{8\pi}}\end{aligned}$$

Hence the normalised version of the wavefunctions are:

$$\begin{aligned}Y_{00} &= \frac{1}{\sqrt{4\pi}} \\ Y_{10} &= \sqrt{\frac{3}{4\pi}} \cos \theta \\ Y_{11} &= -\sqrt{\frac{3}{8\pi}} \sin \theta e^{i\phi}\end{aligned}$$

Q3)

To demonstrate orthogonality we consider:

$$\begin{aligned}
 & \int_0^{2\pi} d\phi \int_0^\pi d\theta \sin\theta Y_{10}(\theta, \phi) Y_{11}^*(\theta, \phi) \\
 &= \int_0^{2\pi} d\phi \int_0^\pi d\theta \sin\theta \left( \sqrt{\frac{3}{4\pi}} \cos\theta \right) \left( -\sqrt{\frac{3}{8\pi}} \sin\theta e^{-i\phi} \right) \\
 &= \frac{3}{\sqrt{32\pi}} \left( \int_0^{2\pi} e^{-i\phi} d\phi \right) \left( \int_0^\pi d\theta \sin^2\theta \cos\theta \right) \\
 &= \frac{3}{\sqrt{32\pi}} \left[ \frac{e^{-i\phi}}{-i} \right]_{2\pi}^0 \left[ \frac{\sin^3\theta}{3} \right]_0^\pi \\
 &= \frac{3}{\sqrt{32\pi}} \left( \frac{1-1}{-i} \right) \left( \frac{0-0}{3} \right) \\
 &= 0
 \end{aligned}$$

Q4) Using the spherical harmonic forms given at the end of Q2, we can express this wavefunction as a sum over  $Y_{00}$ ,  $Y_{10}$  and  $Y_{11}$ :

$$\begin{aligned}
 \psi(\theta, \phi) &= \frac{1}{\sqrt{4\pi}} \left[ \frac{1}{\sqrt{2}} + \cos\theta + \frac{1}{2} \sin\theta e^{i\phi} \right] \\
 &= \frac{1}{\sqrt{2}} \frac{1}{\sqrt{4\pi}} + \frac{1}{\sqrt{4\pi}} \cos\theta + \frac{1}{2} \frac{1}{\sqrt{4\pi}} \sin\theta e^{i\phi} \\
 &= \frac{1}{\sqrt{2}} Y_{00} + \frac{1}{\sqrt{3}} Y_{10} - \frac{1}{\sqrt{6}} Y_{11}
 \end{aligned}$$

The modulus squared of each coefficient gives the probability associated with measuring corresponding eigenvalues. These probabilities are:

| Term in expansion | $(l, m)$       | Probability |
|-------------------|----------------|-------------|
| $Y_{00}$          | $l = 0, m = 0$ | $1/2$       |
| $Y_{10}$          | $l = 1, m = 0$ | $1/3$       |
| $Y_{11}$          | $l = 1, m = 1$ | $1/6$       |

The possible results of measuring  $L^2$  are:

- $l = 0$  or  $L^2 = 0$ , with probability  $\frac{1}{2}$
- $l = 1$  or  $L^2 = 2\hbar^2$ , with probability  $\frac{1}{3} + \frac{1}{6} = \frac{1}{2}$

The possible results of measuring  $L_z$  are:

- $m = 0$  or  $L_z = 0$ , with probability  $\frac{1}{2} + \frac{1}{3} = \frac{5}{6}$
- $m = 1$  or  $L_z = \hbar$ , with probability  $\frac{1}{6}$