Quantum Mechanics Week 3: Class Activities solutions

Q1) In the region |x| < L the potential is V(x) = 0 and the Schrödinger equation takes the form:

$$-\frac{\hbar^2}{2m}\frac{d^2\psi}{dx^2} = E\psi$$

Trying the solution $\psi = A \cos kx$ we find:

$$-\frac{\hbar^2}{2m} \cdot -k^2 A \cos kx = E \cdot A \cos kx$$

which is hence a solution with $k^2 = 2mE/\hbar^2$.

Q2) In the region |x| > L the potential is $V(x) = V_0$ and the Schrödinger equation takes the form:

$$-\frac{\hbar^2}{2m}\frac{d^2\psi}{dx^2} + V_0\psi = E\psi$$

Trying the solution $\psi = Be^{\pm lx}$ we find:

$$-\frac{\hbar^2}{2m} \cdot l^2 B e^{\pm lx} + V_0 \cdot B e^{\pm lx} = E \cdot B e^{\pm lx}$$

which is hence a solution with $l^2 = 2m(V_0 - E)/\hbar^2$.

Q3) The continuity of $\psi(x)$ at x = L implies:

$$A\cos kL = Be^{-lL}$$

where we have taken the decaying form of the exponential solution to ensure $\psi \to 0$ as $x \to \infty$, such that the wavefunction can be normalised. We note that the continuity of $\psi(x)$ at x = -L results in the same equation.

The continuity of $\frac{d\psi}{dx}$ at x = L implies:

$$-kA\sin kL = -lBe^{-lL}$$

Again, the continuity of $\frac{d\psi}{dx}$ at x = -L results in the same equation.

Dividing these two equations we find:

$$\tan kL = \frac{l}{k}$$

Now using the expressions for *k* and *l* from above:

$$\tan\sqrt{\frac{2mEL^2}{\hbar^2}} = \sqrt{\frac{V_0}{E}} - 1$$

Q4) The wavefunction corresponding to the lowest energy state will consist of a cosine solution in the central region with the longest possible wavelength, and decaying exponentials on either side (image credit from http://hyperphysics.phy-astr.gsu.edu/hbase/quantum/pfbox.html):



Q5) The continuity of $\psi(x)$ at x = 0 implies:

$$\begin{split} \Psi_I(0,t) + \Psi_R(0,t) &= \Psi_B(0,t) \\ \to e^{-i\omega t} + Re^{-i\omega t} &= (A+B) e^{-i\omega t} \\ \to 1+R &= A+B \quad [\text{Eq. 1}] \end{split}$$

The continuity of $\psi(x)$ at x = L implies:

$$\Psi_B(x,t) = \Psi_T(x,t)$$

$$\rightarrow (Ae^{kL} + Be^{-kL}) e^{-i\omega t} = Te^{i(kL-\omega t)}$$

$$\rightarrow Te^{ikL} = Ae^{kL} + Be^{-kL} \qquad [Eq. 2]$$

where we have used l = k. The continuity of $\frac{\partial \psi}{\partial x}$ at x = 0 implies:

$$\frac{\partial \Psi_{I}(0,t)}{\partial x} + \frac{\partial \Psi_{R}(0,t)}{\partial x} = \frac{\partial \Psi_{B}(0,t)}{\partial x}$$
$$\rightarrow ike^{-i\omega t} - ik Re^{-i\omega t} = (kA - kB) e^{-i\omega t}$$
$$\rightarrow (1 - R)i = A - B \qquad [Eq. 3]$$

The continuity of $\frac{\partial \psi}{\partial x}$ at x = L implies:

$$\frac{\partial \Psi_B(0,t)}{\partial x} = \frac{\partial \Psi_T(0,t)}{\partial x}$$

$$\rightarrow l(Ae^{kL} - Be^{-kL}) e^{-i\omega t} = ikTe^{i(kL-\omega t)}$$

$$\rightarrow i Te^{ikL} = Ae^{kL} - Be^{-kL} \qquad [Eq. 4]$$

Q6) We can now find the unknown coefficients from Eq.1-4. First we consider:

$$[Eq. 2] + [Eq. 4] \rightarrow (1 + i) Te^{ikL} = 2Ae^{kL}$$
 [Eq. 5]
 $[Eq. 2] - [Eq. 4] \rightarrow (1 - i) Te^{ikL} = 2Be^{-kL}$ [Eq. 6]

Dividing [Eq.5] by [Eq.6] gives:

$$\frac{B}{A} = \left(\frac{1-i}{1+i}\right) e^{2kL} \qquad [\text{Eq. 7}]$$

Re-arranging [Eq.5] gives:

$$\frac{T}{A} = \frac{2 e^{kL} e^{-ikL}}{1+i}$$
 [Eq. 8]

Now we determine *A* using the other 2 equations:

$$[Eq. 1] - [Eq. 3] \times i \to 2 = (1 - i)A + (1 + i)B$$

$$\to 2 = (1 - i)A + (1 + i)\left(\frac{1 - i}{1 + i}\right)e^{2kL}A$$

$$\to 2 = A(1 - i)(1 + e^{2kL})$$

$$\to A = \frac{2}{(1 - i)(1 + e^{2kL})}$$

Now we find *T* using [Eq.8]:

$$T = \frac{2 e^{kL} e^{-ikL}}{1+i} \cdot \frac{2}{(1-i)(1+e^{2kL})} = \frac{2 e^{-ikL}}{e^{kL} + e^{-kL}}$$

which leads to:

$$|T|^2 = \frac{4}{(e^{kL} + e^{-kL})^2}$$