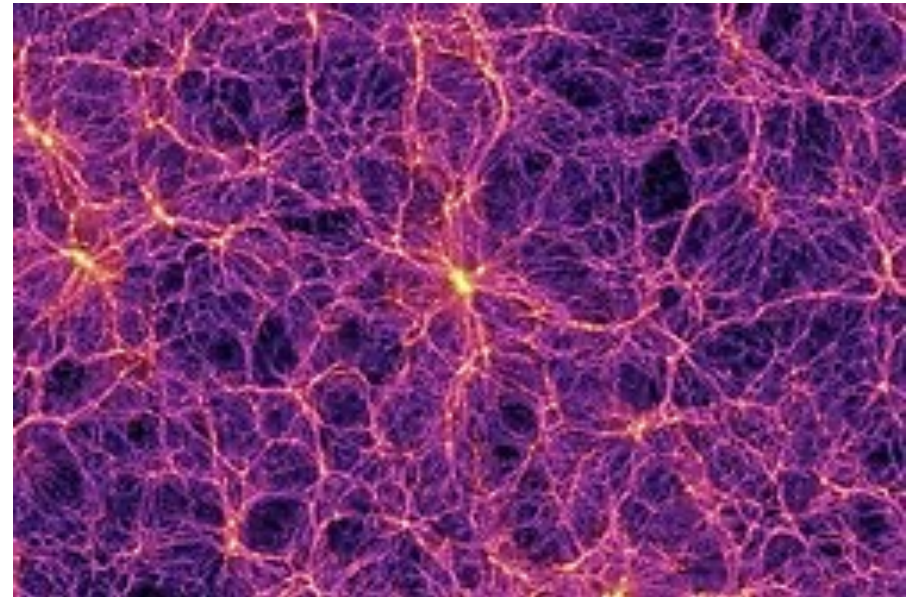
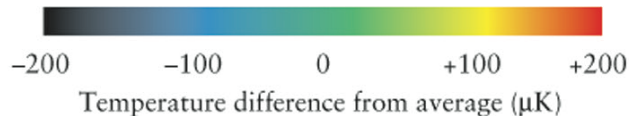
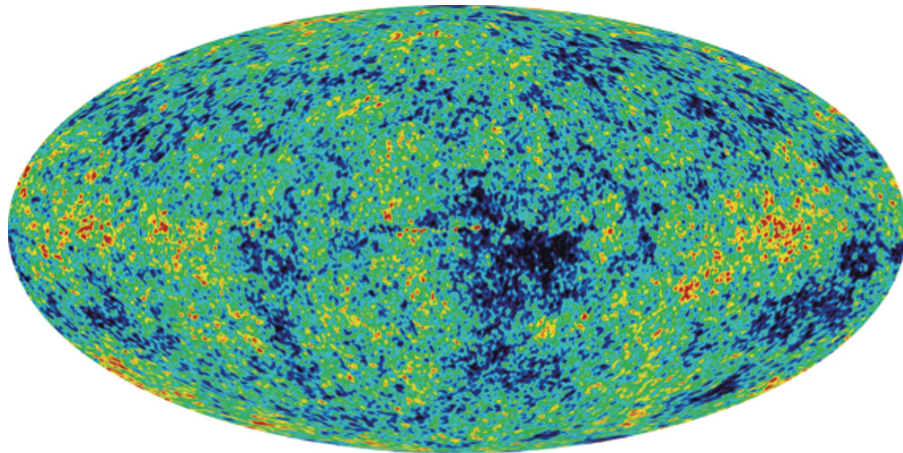


# Honours Cosmology Week 5: Cosmic Structure

*This week we will study the processes by which the large-scale structures in the Universe are created and gravitationally grow with time*



# Cosmic structure

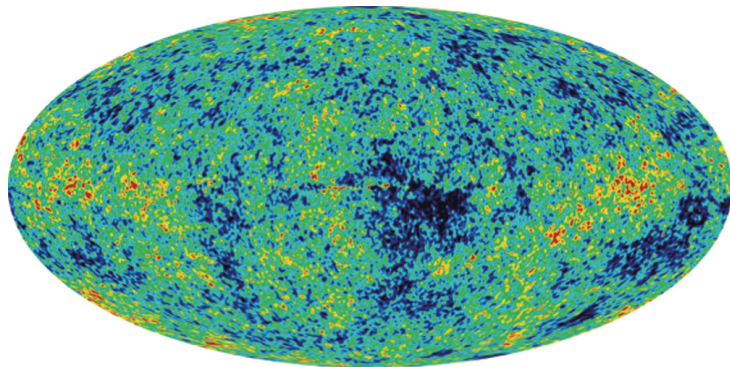
At the end of this week you should be able to ...

- ... motivate the theory of **inflation**, which might seed the original density fluctuations in the Universe
- ... understand how inflation can address the **flatness problem** and **horizon problem**
- ... perform calculations for the **gravitational collapse time** of density clumps, and relate them to real cosmic objects
- ... describe that structure grows in a **hierarchical fashion**
- ... recognise observable effects of the growth of structure, such as **gravitational lensing** and **redshift-space distortions**

# Cosmic structure

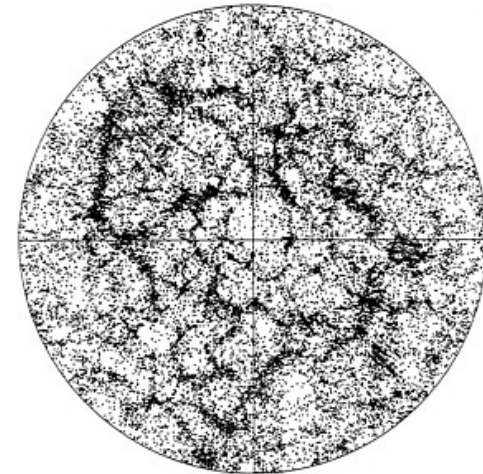
- Primordial clumps of matter in the Universe are *gravitationally amplified* to form today's galaxies

CMB fluctuations with  $\Delta T/T \sim 10^{-5}$



-200    -100    0    +100    +200  
Temperature difference from average ( $\mu\text{K}$ )

Distribution of galaxies with  
 $\rho_{\text{galaxy}}/\rho_{\text{crit}} \sim 10^6$



- In this week's class, we'll study the physics and observational consequences of this process

# Flatness problem

- We'll start with the question: *what seeded the original fluctuations?* This question is connected with two other issues ...
- First, let's consider the *curvature* of the Universe. The Universe today is measured to be flat,  $|\Omega_K| < 0.01$ . What is the flatness of the Universe at other times?
- Density parameters *as seen by observers at different times* are given by, using  $\rho_{\text{crit}}(t) = 3H(t)^2/8\pi G$ ,

$$\frac{\Omega(t)}{\Omega(t_0)} = \frac{\rho(t)/\rho_{\text{crit}}(t)}{\rho(t_0)/\rho_{\text{crit}}(t_0)} = \frac{\rho(t)}{\rho_0} \frac{H_0^2}{H^2}$$

This comes from the Friedmann equation

This depends on the component

# Flatness problem

- Assuming the Universe is matter-dominated,  $\frac{H^2}{H_0^2} \approx 1/a^3$ , and for curvature,  $\frac{\rho_K(t)}{\rho_{K,0}} = \frac{1}{a^2}$ . Hence:  $\Omega_K(t) = \Omega_{K,0} \cdot a$
- If the Universe is flat today ( $\Omega_{K,0} = 0$ ), it will always be flat ( $\Omega_K(t) = 0$ ). *A small  $\Omega_K$  in the past increases  $\propto a$ !*
- If  $|\Omega_K| < 0.01$  today, then  $|\Omega_K| < 10^{-12}$  in the first moments of the Universe ( $a \sim 10^{-10}$ )
- What makes the early Universe flat to this incredible precision? This question is known as the *flatness problem*

# Horizon problem

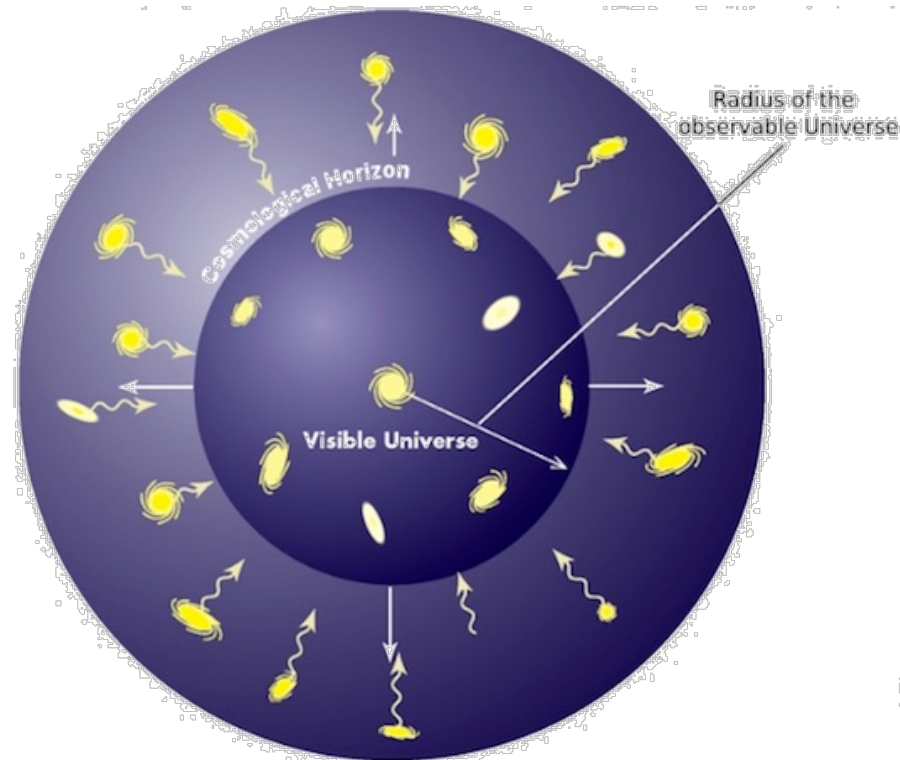
- The *particle horizon* is the maximum distance,  $D_H$ , that a particle can travel between the Big Bang and a given time, and hence gives the *maximum range of causal connection*

- Physical distance between points  $A$  and  $B$ , measured at a simultaneous coordinate time =  $a(t) \int_A^B \frac{dr}{\sqrt{1-Kr^2}}$

- Suppose these points are linked by a light ray, which travels with

$$ds = 0 \rightarrow c dt' = a(t') \frac{dr}{\sqrt{1-Kr^2}}$$

- Hence,  $D_H(t) = c \cdot a(t) \int_0^t \frac{dt'}{a(t')}$



# Horizon problem

- Note that the horizon distance is **not** equal to  $D_H(t) = ct$ , since the Universe is expanding during this time, so points  $A$  and  $B$  end up **further apart!**
- Assuming the Universe has  $\Omega_m = 1$ , we have  $a(t) = \left(\frac{3H_0 t}{2}\right)^{2/3}$  and  $D_H(t) = c \cdot a(t) \int_0^t \frac{dt'}{a(t')}$  evaluates to:

$$D_H(t) = 3ct = \frac{2c}{H_0} a^{3/2}$$

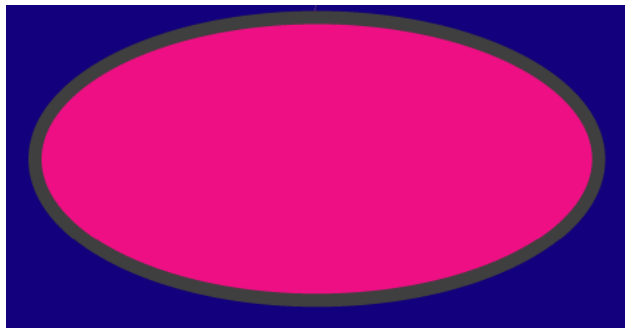
- Assuming  $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$ , the particle horizon today has size  $D_H = 8.6 \text{ Gpc}$ , however it's much smaller in the early Universe!

# Horizon problem

- The CMB is produced at  $a \approx 10^{-3}$  and its temperature is uniform across the sky. The range of *causal connection* must therefore correspond to the size of the sky,  $\Delta\theta \sim 180^\circ$  !
- However, using the metric to project  $D_H$  at  $a = 10^{-3}$  to an angular separation on the sky we find,

$$\Delta\theta = \frac{D_H}{r a} = \frac{\frac{2c}{H_0} a^{1/2}}{\frac{2c}{H_0} (1-\sqrt{a})} = \sqrt{10^{-3}} = 0.03 \text{ rad} = 1.8^\circ$$

The CMB temperature varies by only 1 part in  $10^5$  across the sky



However, the causal connection distance is only a tiny angle! This is known as the **horizon problem**

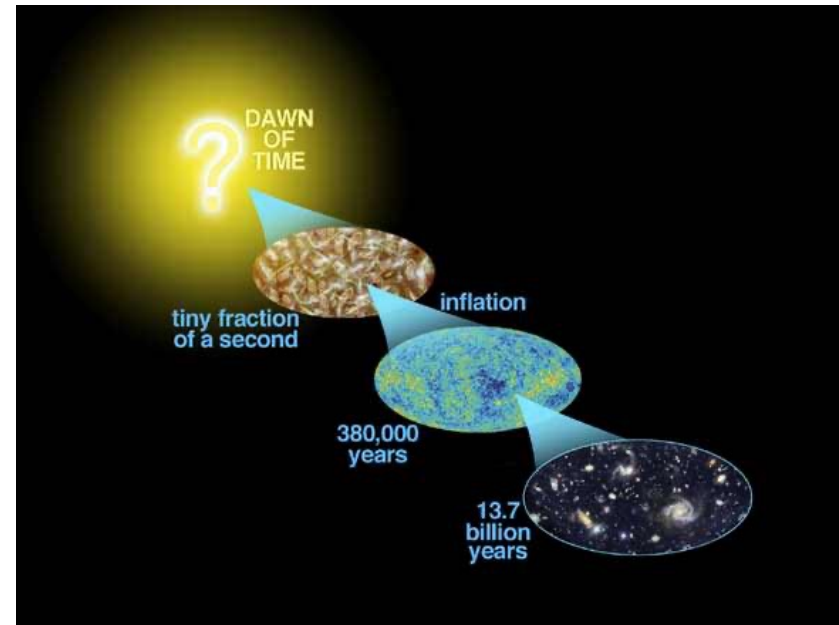
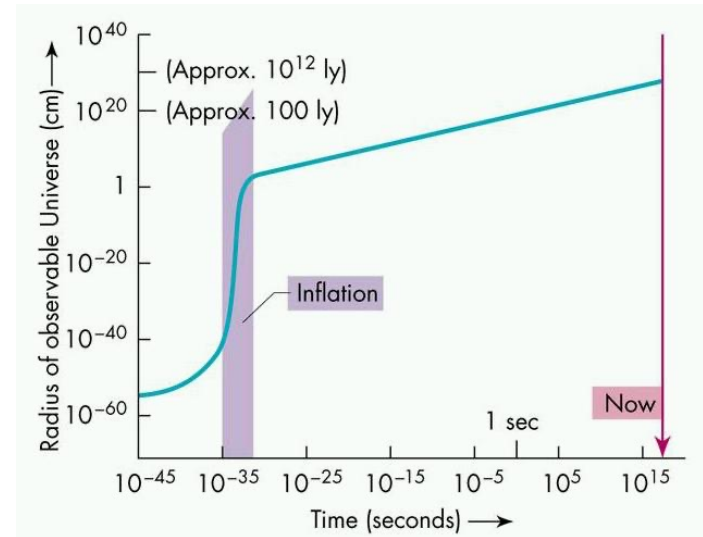


# Inflation

- A simple and elegant solution to both of these problems is if the early Universe *had a brief period of accelerated expansion*
- **The flatness problem:** accelerated expansion drives  $\Omega_K$  closer to zero with time [e.g. repeating the calculation on slide 5 for dark energy:  $\frac{H^2}{H_0^2} = 1$ , hence  $\Omega_K(t) = \Omega_{K,0}/a^2$  which decreases with time!]
- **The horizon problem:** accelerated expansion increases the particle horizon by allowing nearby regions to expand to greater separations, bringing the CMB sky into causal contact
- This period of early accelerated expansion is called *inflation*

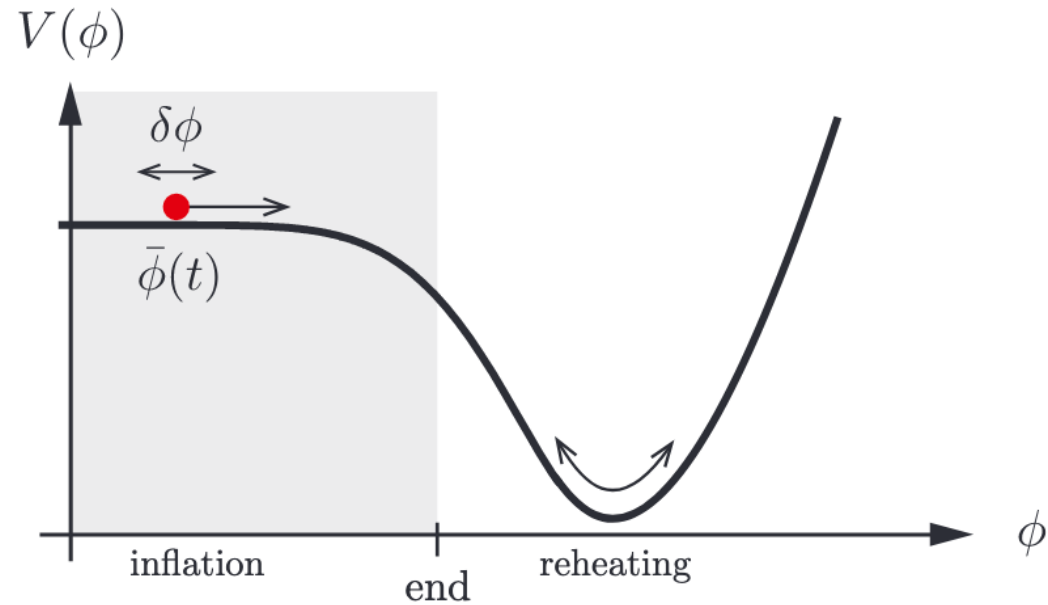
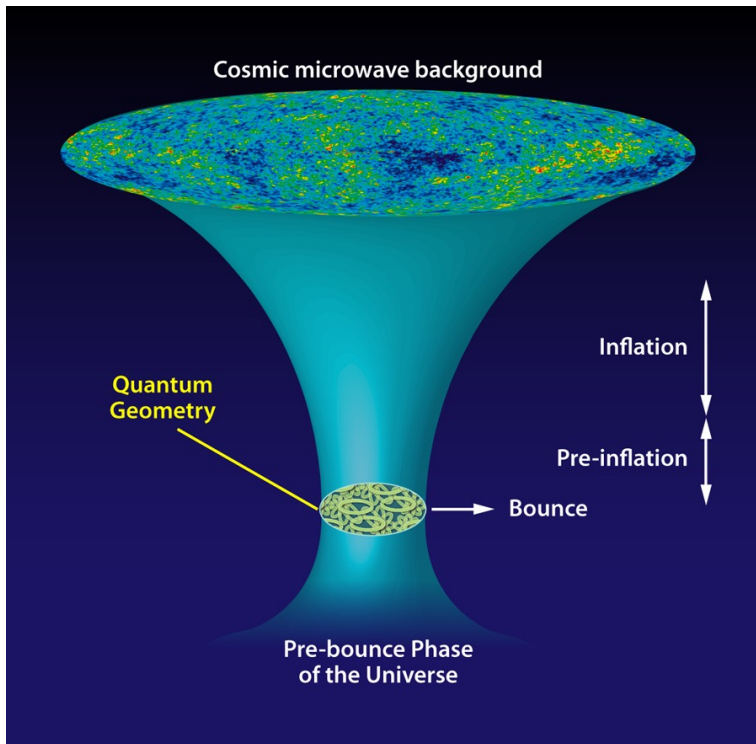
# Inflation

- Inflation is thought to occur within the first fraction of a second of the Universe, and expand its size by  $\sim 10^{40}$
- Inflation can also seed the density fluctuations that are seen in the Universe by *amplifying quantum fluctuations to a macroscopic scale*



# Inflation

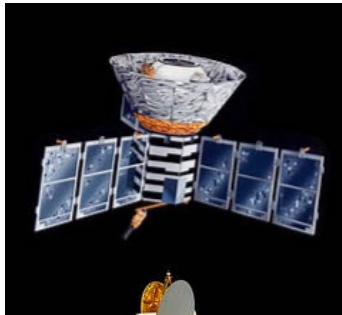
- The ultimate cause of inflation is currently unknown!
- It would be intriguing if it were related to the other period of accelerating expansion: dark energy



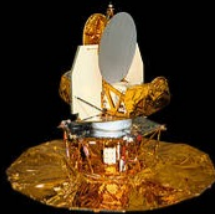
# CMB fluctuations

- A series of satellite experiments have mapped out the tiny fluctuations in the CMB with increasing sensitivity and resolution

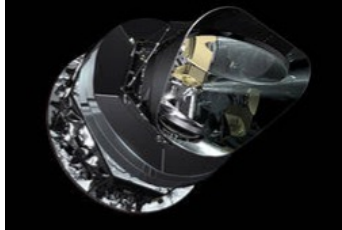
**COBE**



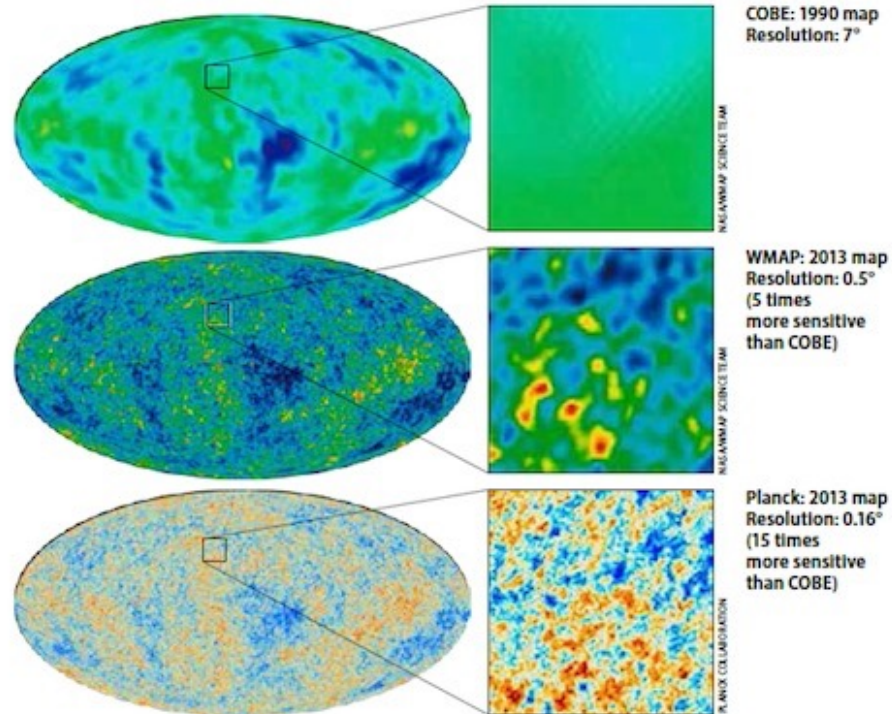
**WMAP**



**Planck**

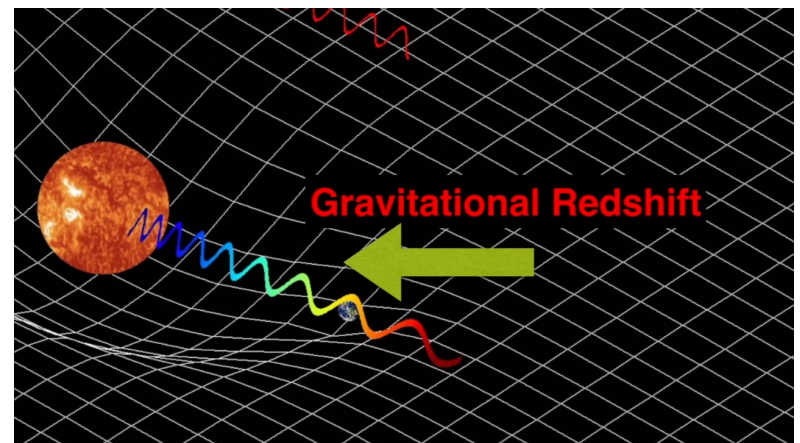
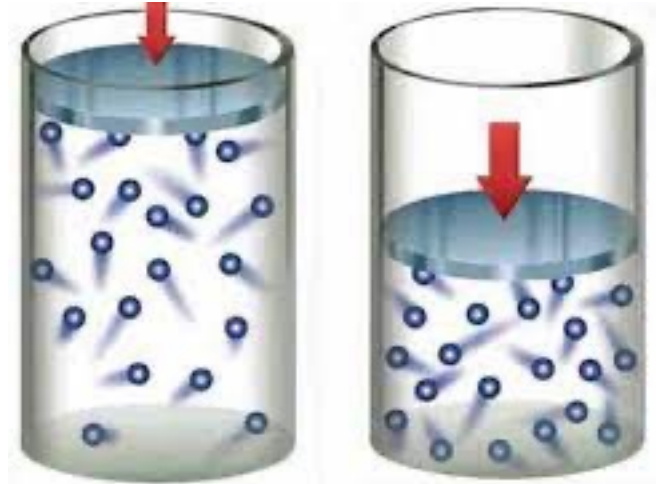


## CLOSE-UP VIEWS OF THE CMB



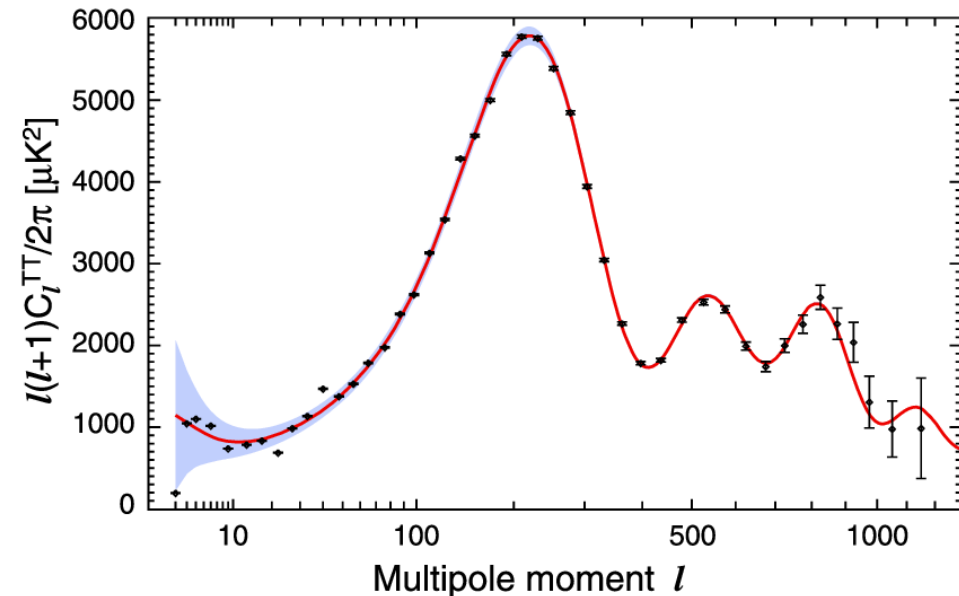
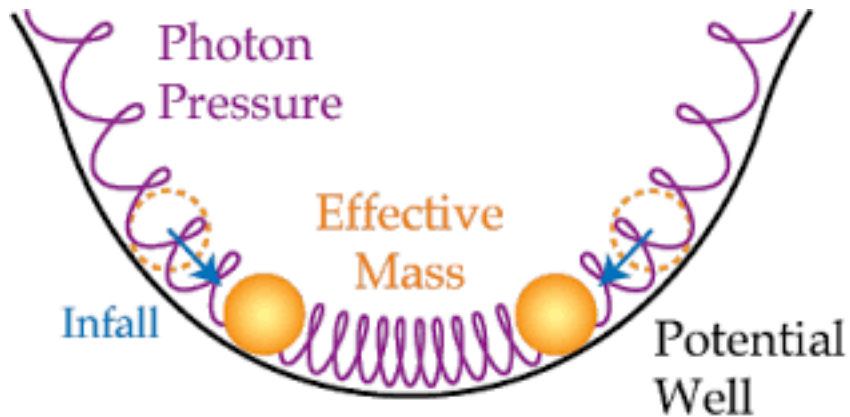
# CMB fluctuations

- Two main effects transform the matter fluctuations to variations in temperature
- *Adiabatic compression*: a gravitating clump of matter squeezes the gas, implying a **higher temperature**
- *Gravitational redshift*: light originating from a gravitating clump loses energy as it escapes, implying a **lower temperature**



# CMB fluctuations

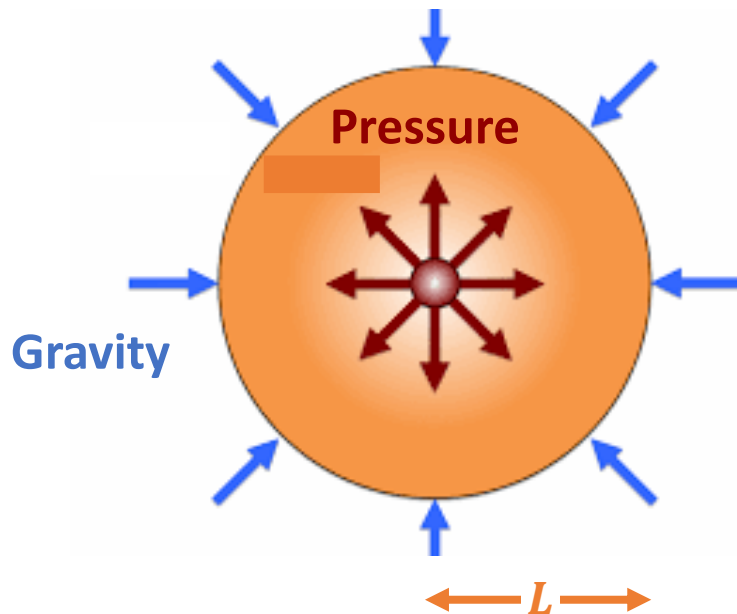
- These competing forces create compression waves or **acoustic waves** in the plasma, which may be observed today in the CMB temperature spectrum (we'll look at this in more detail next week)



# Spherical collapse

- The balance between the inward gravitational force and outward pressure force determines when a clump collapses under its own weight

Clump of  $N$  particles of mass  $m$   
with density  $\rho$  and temperature  $T$



Here is an approximate calculation of the minimum collapse size or **Jeans length**:

Number of particles  $N \sim \rho L^3 / m$

Gravitational potential energy of the clump  $E_G \sim -\frac{G(Nm)^2}{L}$

Thermal energy of the clump  $E_T \sim Nk_B T$

Gravity dominates if  $E_G + E_T < 0$

i.e. if  $L > \sqrt{\frac{k_B T}{G \rho m}} = \text{Jeans length } L_J$

# Spherical collapse

- We can learn from the calculation on the previous slide that **large regions** ( $L > L_J$ ) **collapse** under gravity whilst **small regions** ( $L < L_J$ ) **oscillate** between pressure and gravity
- The **collapse time**  $t_G$  can be estimated assuming a constant gravitational acceleration  $\frac{GM}{L^2} \sim \frac{L}{t_G^2} \rightarrow t_G \sim \frac{1}{\sqrt{G\rho}}$
- The minimum mass for collapse, or **Jeans mass**  $M_J$ , can be estimated as  $M_J \sim \rho L_J^3 = \rho \left( \frac{k_B T}{G\rho m} \right)^{3/2} \propto T^{3/2} \rho^{-1/2}$
- We learn that the mass needed for collapse reduces with **lower temperature** or **higher density**



# Spherical collapse

- We can calculate the *gravitational collapse time* of some typical objects in the Universe:

## Globular cluster

$$t_G \sim 2 \text{ Myr}$$



Can form very early in the history of the Universe

## Galaxy

$$t_G \sim 300 \text{ Myr}$$



In place by  $z \sim 10$

## Galaxy cluster

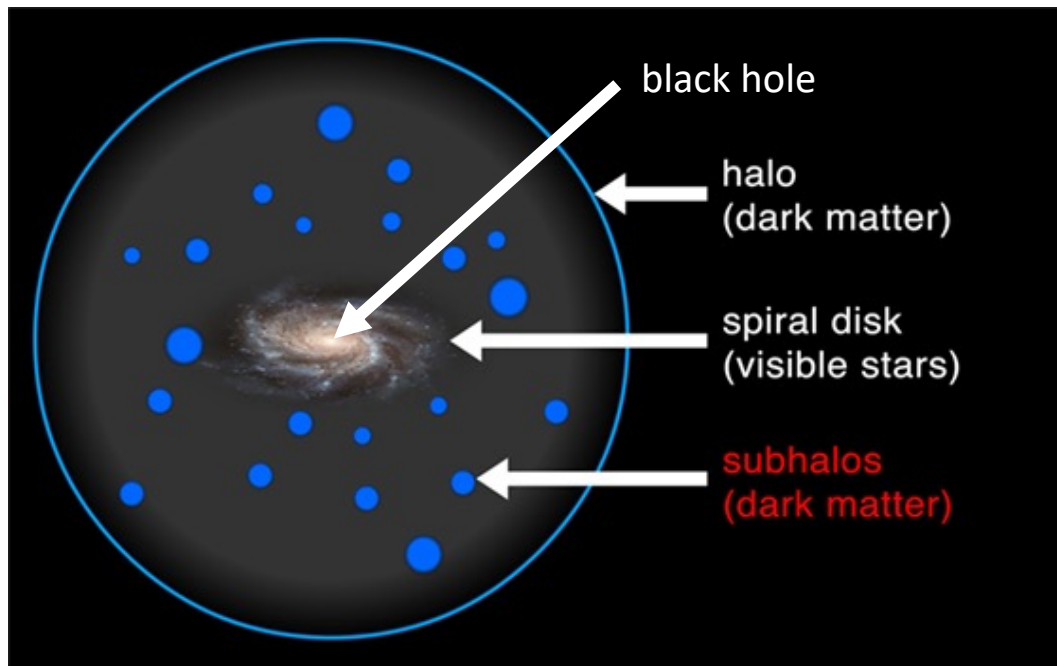
$$t_G \sim 7 \text{ Gyr}$$



In place by  $z \sim 1$

# Hierarchical structure formation

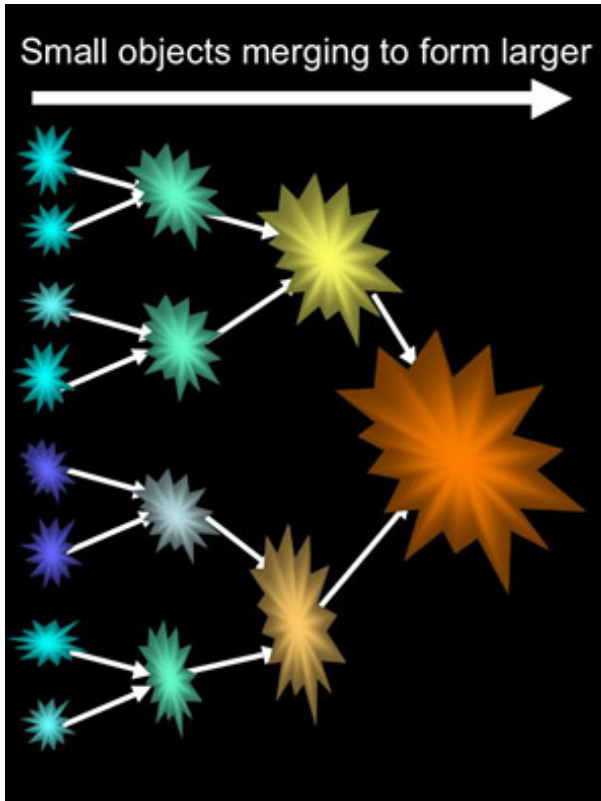
- Dark matter clumps collapse early into “*dark matter halos*” and form the gravitating structure of the Universe
- Baryons fall into these haloes and form galaxies, which contain *supermassive black holes* at their centre



- Dark matter halos have a characteristic “NFW” (Navarro-Frenk-White) radial mass profile
- Dark matter halos contain substructure which may host satellite galaxies

# Hierarchical structure formation

- The large-scale distribution of galaxies forms in a *hierarchical fashion* – smaller structures form first, and agglomerate into larger ones!



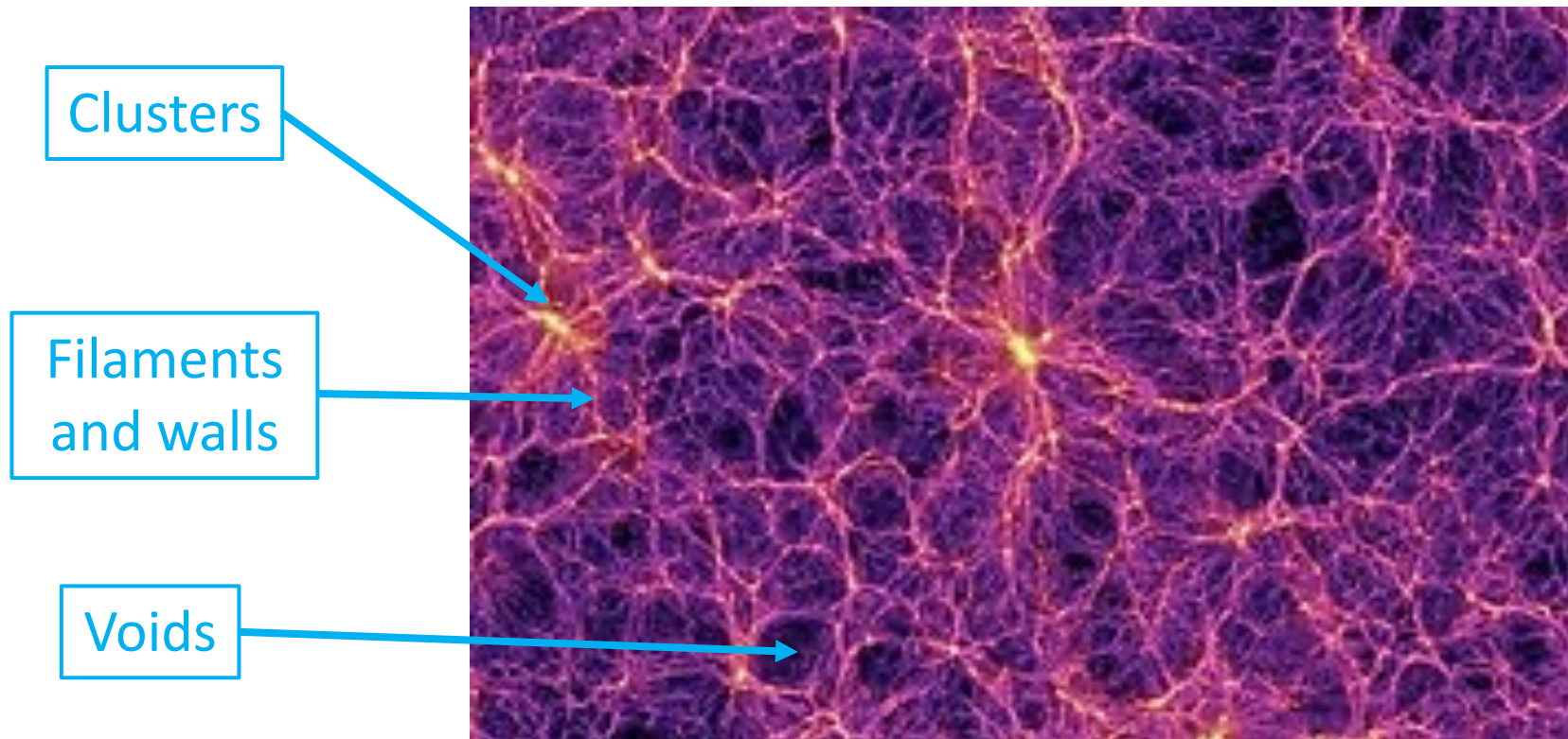
We can observe galaxies in the process of merging!



# Hierarchical structure formation

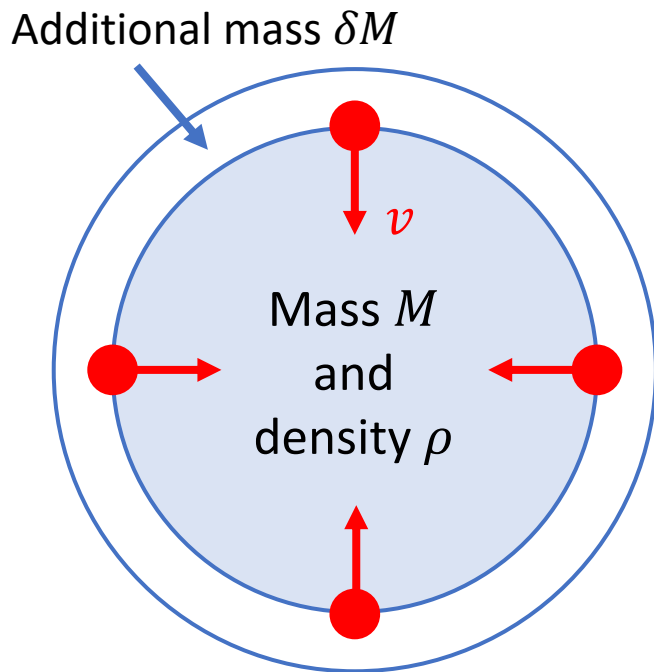
- A *sponge-like topological structure* of galaxies results, featuring clusters, walls, filaments and voids!

Simulation of the matter distribution within a large volume



# Infall velocity

- The gravitational collapse of a structure will create an *infall velocity* in the surrounding matter



We can relate the increase in mass  $\delta M$  in time  $\delta t$  to the infall velocity  $v$  by the **continuity equation**. The additional mass is in a slice of thickness  $v \times \delta t$ :

$$\delta M = (v \delta t) \cdot 4\pi r^2 \cdot \rho$$

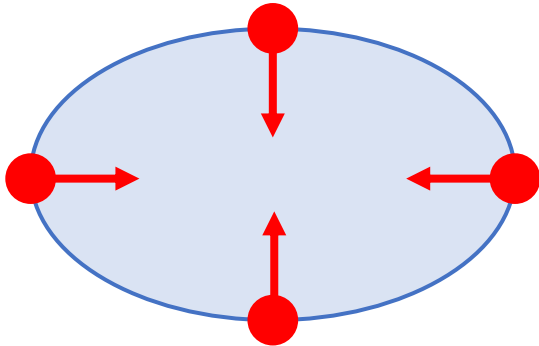
Since  $M = \frac{4}{3}\pi r^3 \rho$ , we have

$$v = \frac{1}{M} \frac{dM}{dt} \cdot \frac{r}{3}$$

# Redshift-space distortions

- This infall velocity is observable because it systematically changes the galaxy redshifts in an affect known as *redshift-space distortions*

This galaxy is moving towards the observer so is **blueshifted**

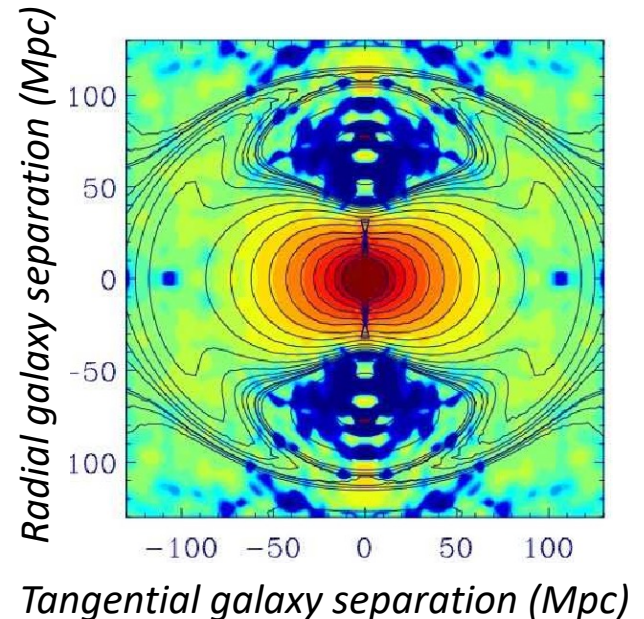


This galaxy is moving away from the observer so is **redshifted**



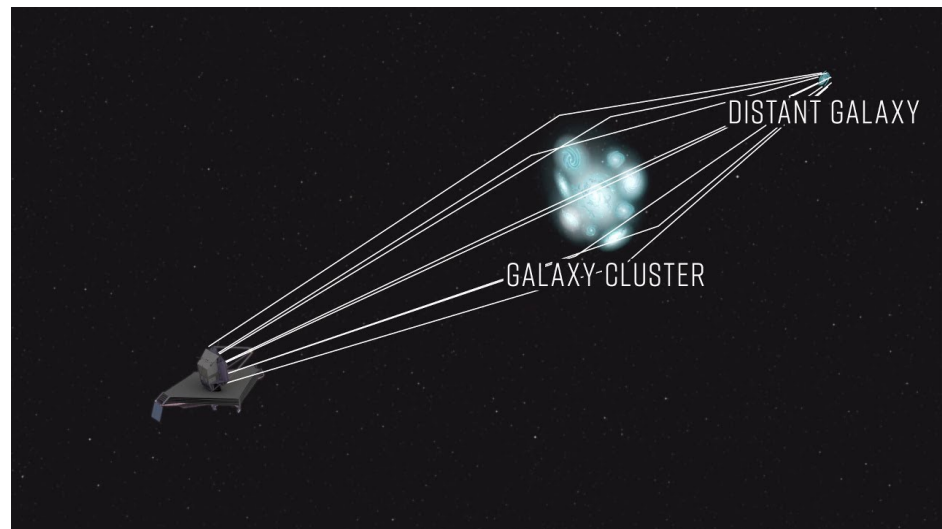
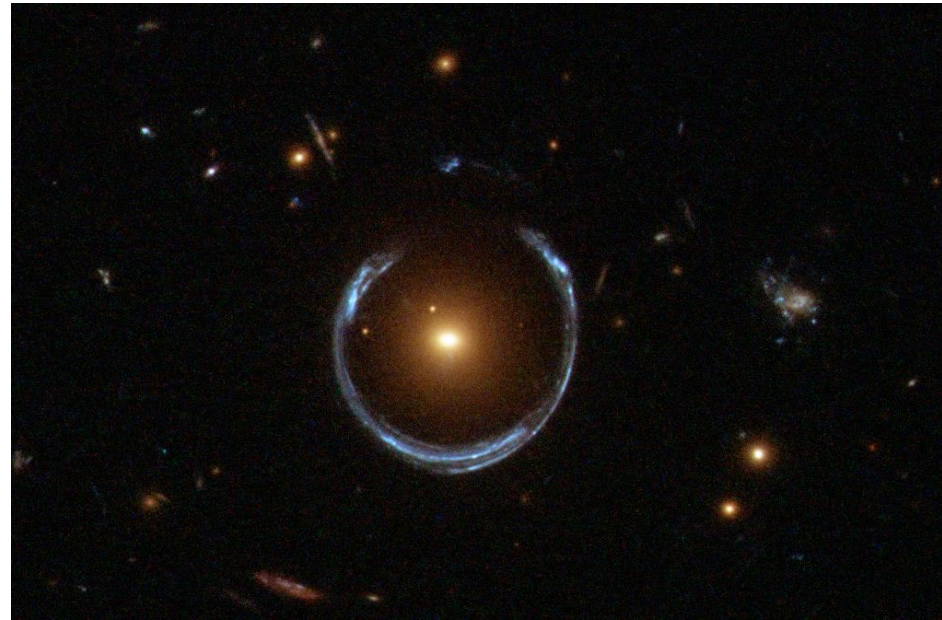
The “average structure” appears squashed when mapped out with galaxy redshifts

This is a real observation made by stacking up structures in redshift-space:



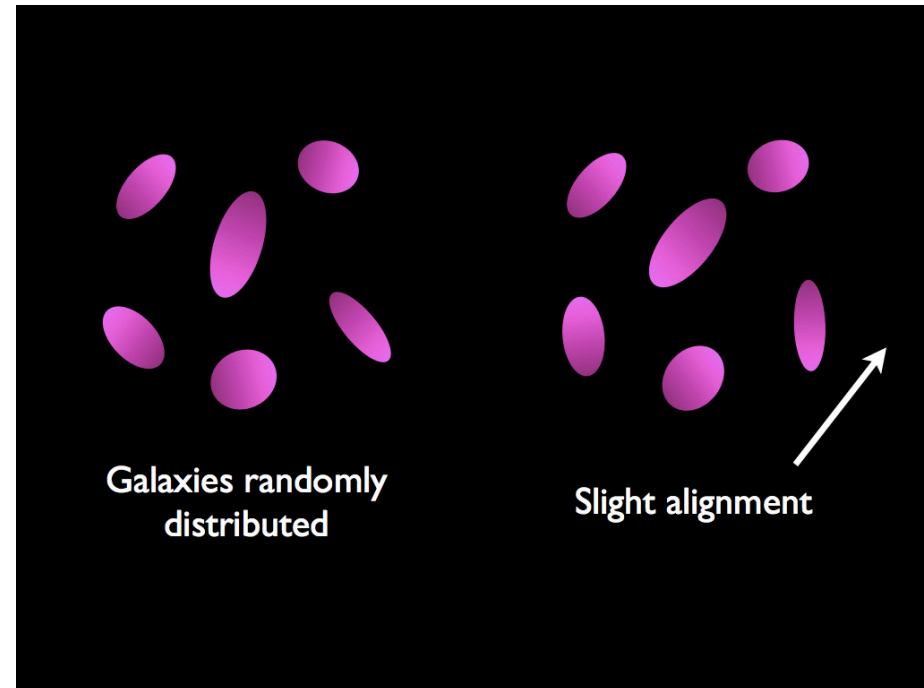
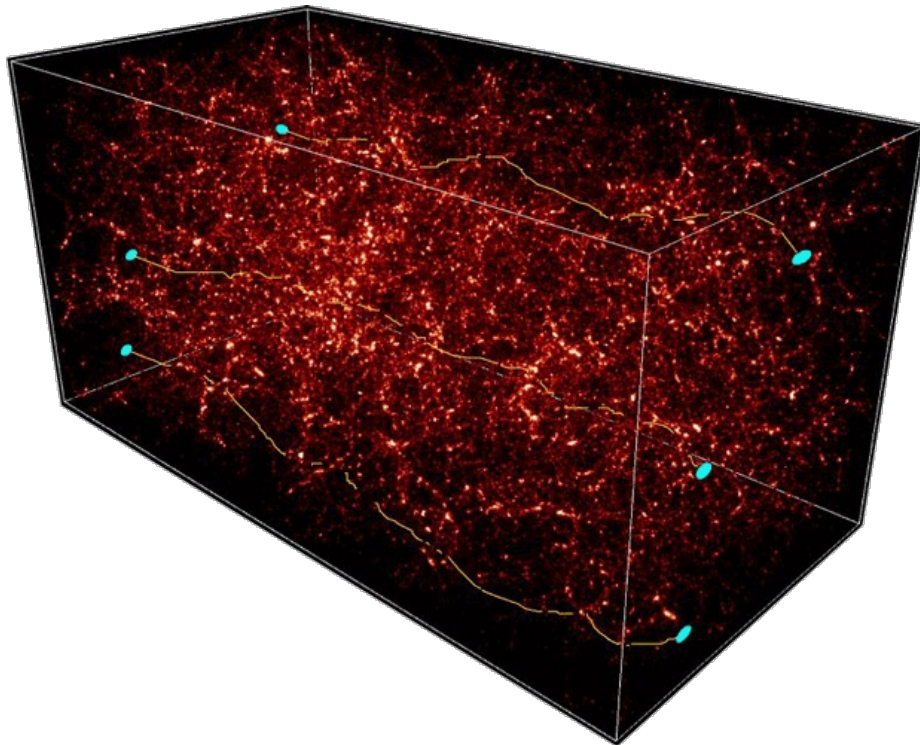
# Gravitational lensing

- Massive clumps cause *gravitational lensing* of light from background galaxies, which can be observed through characteristic arcs
- The space-time around the masses is *distorted* according to General Relativity, bending the paths of the light rays



# Gravitational lensing

- An important technique called *cosmic shear* uses the small lensing distortions in background galaxy shapes to trace out the large-scale structure





# Key take-aways

- According to our current theories, the density fluctuations in the Universe were created by a very early period of rapid expansion known as **inflation**
- Inflation is an attractive theory because it can solve both the **flatness problem** and **horizon problem**
- The density fluctuations can be traced as **temperature fluctuations in the CMB** and by the **growth of galaxies**
- Cosmic structure grows in a **hierarchical manner**, with smaller objects forming early and merging into larger objects
- The growth of structure may be measured by effects such as **gravitational lensing** and **redshift-space distortions**