Honours Cosmology Week 2: Measuring the Universe

This week we'll explore the space-time metric of the Universe, and how it may be used to describe cosmic distances and light travel


## Measuring the Universe

At the end of this week you should be able to ...

- ... understand how metrics may be used to compute distances in a given coordinate system and topology
- ... understand the form of the Friedmann-Robertson-Walker (FRW) metric of a homogeneous \& isotropic Universe
- ... use this metric to write down the equation for light propagation in an expanding Universe
- ... describe different definitions of cosmological distance
- ... determine the age of the Universe given the scale factor


## Measuring the Universe

- This week we'll learn about the equations for mapping out the Universe - that is, how to relate distances, redshifts, angles and light travel time



## Introducing metrics

- Let's first consider how to measure distances in a 2D Cartesian space with coordinates ( $x, y$ )

- We break the arc into small pieces: $d s^{2}=d x^{2}+d y^{2}$
- Integrating along the curve: $s=\int_{A}^{B} \sqrt{\left(\frac{d x}{d \lambda}\right)^{2}+\left(\frac{d y}{d \lambda}\right)^{2}}$ $d \lambda$
- The relation between the line element $d s$ and coordinate intervals $(d x, d y)$ is called the metric of the space


## Introducing metrics

- The metric depends on the chosen co-ordinate system

For 3D Cartesian coordinates, we have:

$$
d s^{2}=d x^{2}+d y^{2}+d z^{2}
$$

Describing exactly the same space with spherical polar co-ordinates:

$$
d s^{2}=d r^{2}+r^{2}\left(d \theta^{2}+\sin ^{2} \theta d \phi^{2}\right)
$$



For example, to calculate the distance to circumnavigate the world from Melbourne on a line of constant latitude $\theta_{0}: d r=0$ and $d \theta=0$, so

$$
s=\int d s=r_{\text {Earth }} \sin \theta_{0} \int_{0}^{2 \pi} d \phi=2 \pi r_{\text {Earth }} \sin \theta_{0}
$$

## Introducing metrics

- The metric also depends on the intrinsic topology of the space
- To describe the constant curvature surface of a unit sphere, we could use spherical polar co-ordinates with $r=1$ :

$$
d s^{2}=d \theta^{2}+\sin ^{2} \theta d \phi^{2}
$$



- A 2D flat space would be described by $d s^{2}=d \theta^{2}+d \phi^{2}$ : in this case, the metrics are different but so are the intrinsic topologies (curved space vs. flat space)


## Curved spaces

## - How do curved spaces differ from flat spaces?



- In a curved space: parallel lines converge or diverge
- The circumference of a circle $\neq 2 \pi r$
- The sum of the angles of a triangle $\neq 180^{\circ}$
- The curvature of a space can be determined by local observers without needing to see the space from outside


## Curved spaces

- How do we generalise the metric of a constant curvature space from 2D, as above, to 3D?
- A constant curvature 2D surface embedded in a 3D Euclidean space satisfies the equation: $x^{2}+y^{2}+z^{2}=1 / K \quad(K=$ curvature parameter $)$
- For a constant curvature 3D surface embedded in a 4D Euclidean space: $x^{2}+y^{2}+z^{2}+w^{2}=1 / K \quad$ (adding the extra co-ordinate $w$ )
- In the 4D Euclidean space: $d s^{2}=d x^{2}+d y^{2}+d z^{2}+d w^{2}$
- We now transform $(x, y, z)$ to spherical polar co-ordinates $(r, \theta, \phi)$
- From above: $w^{2}=1 / K-r^{2}$, hence $d w^{2}=\frac{r^{2} d r^{2}}{1 / K-r^{2}}$
- Putting it all together: $d s^{2}=\frac{d r^{2}}{1-K r^{2}}+r^{2}\left(d \theta^{2}+\sin ^{2} \theta d \phi^{2}\right)$


## Introducing General Relativity

- The Universe is described by Einstein's theory of General Relativity, summarised in two phrases:

Space-time tells matter how to move

Matter tells space-time how to curve


- We won't study GR in detail in this module, but we'll use a key aspect of it: the space-time metric


## The metric of the Universe

- In relativity, the metric describes separations in "space-time" rather than just "space"
- We already know the space-time interval in special relativity (invariant for observers in all inertial frames):

$$
d s^{2}=-c^{2} d t^{2}+d x^{2}
$$

- We obtain the metric of the Universe by replacing $d x$ with the line element for comoving coordinates of a constant curvature space expanding with the Universe
- We must preserve constant curvature to satisfy the conditions that the Universe is homogeneous and isotropic


## The metric of the Universe

- Combining these results, we obtain the space-time metric of the expanding Universe:

$$
\begin{aligned}
& \qquad d s^{2}=-c^{2} d t^{2}+a(t)^{2}\left[\frac{d r^{2}}{1-K r^{2}}+r^{2}\left(d \theta^{2}+\sin ^{2} \theta d \phi^{2}\right)\right] \\
& \begin{array}{l|l|l}
\text { Same term as in special } \\
\text { relativity, where } t \text { is the } \\
\text { time measured by } \\
\text { fundamental observers }
\end{array} \\
& \hline \begin{array}{l}
\text { Multiplied by the cosmic } \\
\text { scale factor to convert } \\
\text { comoving coordinates to } \\
\text { physical separation }
\end{array}
\end{aligned} \begin{array}{ll}
\begin{array}{l}
\text { Line element of a space } \\
\text { with constant curvature } K \\
\text { in comoving coordinates } \\
(r, \theta, \phi)
\end{array} \\
\hline
\end{array}
$$

- This form is called the Friedmann-RobertsonWalker (FRW) metric, after its discoverers


## Light travel through the Universe

- We can use this spacetime metric to study a light ray travelling through the Universe on a radial path from a galaxy at coordinate $r_{\mathrm{em}}$ to $r=0$

(C) CanStockPhoto.com
- Remember from special relativity, light travels between two points of space-time such that $d s=0$
- Since $d \theta=d \phi=0$, we find: $c d t=-\frac{a(t) d r}{\sqrt{1-K r^{2}}}$

Choosing the negative solution of the square root, since $r$ is decreasing

- Integrating along the path, $\int_{t_{\mathrm{em}}}^{t_{\mathrm{obs}}} \frac{d t}{a(t)}=\frac{1}{c} \int_{0}^{r_{\mathrm{em}}} \frac{d r}{\sqrt{1-K r^{2}}}$


## Light travel through the Universe

- For a second photon emitted slightly later and following the first, we would likewise find: $\int_{t_{\mathrm{em}}+\delta t_{\mathrm{em}}}^{t_{\mathrm{obs}}+\delta t_{\mathrm{obs}}} \frac{d t}{a(t)}=\frac{1}{c} \int_{0}^{r_{\mathrm{em}}} \frac{d r}{\sqrt{1-K r^{2}}}$
- Since the right-hand side of the previous two equations is the same, we can conclude: $\int_{t_{\mathrm{em}}}^{t_{\mathrm{obs}}} \frac{d t}{a(t)}=\int_{t_{\mathrm{em}}+\delta t_{\mathrm{em}}}^{t_{\mathrm{ebs}}+\delta t_{\mathrm{ob}}} \frac{d t}{a(t)}$
- From which we deduce the key result: $\frac{\delta t_{\mathrm{obs}}}{a\left(t_{\mathrm{obs}}\right)}=\frac{\delta t_{\mathrm{em}}}{a\left(t_{\mathrm{em}}\right)}$
- The time between the photons is inversely proportional to the frequency $\omega$ of the light, $\delta t \propto 1 / \omega \propto$ the wavelength $\lambda$
- Hence the light is redshifted: $\lambda_{\mathrm{obs}}=\frac{a\left(t_{\mathrm{obs}}\right)}{a\left(t_{\mathrm{em}}\right)} \lambda_{\mathrm{em}}=(1+z) \lambda_{\mathrm{em}}$


## Distances in expanding space

- How do we describe distances in the expanding Universe, from $r=0$ to a galaxy at coordinate $r$ ?



## Distances in expanding space

- The first approach we might consider is evaluating the physical distance (or "proper distance") between the origin and the galaxy at a chosen time. From the metric:

$$
d D=\frac{a(t)}{\sqrt{1-K r^{2}}} d r \quad(\text { since } d \theta=d \phi=d t=0)
$$

- Integrating between 0 and $r$ we find results depending on $K$ :

$$
\begin{array}{ll}
D=a(t) \sin ^{-1}(r \sqrt{K}) / \sqrt{K} & K>0 \\
D=a(t) r & K=0 \\
D=a(t) \sinh ^{-1}(r \sqrt{-K}) / \sqrt{-K} & K<0
\end{array}
$$

## Distances in expanding space

- However, the physical distance is not a very useful quantity in observational cosmology, since we can't in practice arrange for such a measurement to be made!
- How can we actually measure distances?
- A nice method would be to use the shift in angular position of an object due to the orbit of the Earth, called a parallax distance



## Distances in expanding space

- More useful distance measures are based on the observed size or brightness of distant objects

| Distance |
| :---: |
| measurements |
| based on the |
| size or |
| separation of |
| distant objects |
| produce an |
| angular |
| diameter |
| distance |



Distance measurements based on the brightness of distant objects produce a luminosity distance

## Angular diameter distance

- Angular diameter distance $D_{A}$ is defined using the observed angular size $\Delta \theta$ of an object of known physical width $W$

- The angular diameter distance is defined as: $D_{A}=W / \Delta \theta$
- From the metric: $W=a(t) r \Delta \theta$, where $r=$ object coordinate
- Hence: $D_{A}=a\left(t_{\mathrm{em}}\right) r=\frac{r}{1+z} \quad\left(t_{\mathrm{em}}=\right.$ emission time $)$
- Objects of known physical width $W$ are called standard rulers


## Luminosity distance

- Luminosity distance $D_{L}$ is defined using the flux $f$ received from an object of known luminosity $L$

Luminosity L (known)


Flux $f$ (observed)

- The usual "inverse square law" would give us: $f=\frac{L}{4 \pi D_{L}{ }^{2}}$
- This motivates our definition: $D_{L}=\sqrt{\frac{L}{4 \pi f}}$


## Luminosity distance

- To relate the luminosity distance to the object's coordinate $r$, we have to take two cosmological effects into account:

1. Photons lose energy as their wavelength increases owing to the expansion of the Universe [effect $\propto a\left(t_{\mathrm{em}}\right)$ ]
2. The time interval between arriving photons also increases (see the calculation on slide 13) [effect $\left.\propto a\left(t_{\mathrm{em}}\right)\right]$

- Hence, the flux of energy received is: $f=\frac{L a\left(t_{\mathrm{em}}\right)^{2}}{4 \pi r^{2}}$
- Hence: $D_{L}=\sqrt{\frac{L}{4 \pi f}}=\frac{r}{a\left(t_{\mathrm{em}}\right)}=r(1+z)$
- Objects of known luminosity $L$ are called standard candles


## Volumes

- The volume element in comoving space between coordinates $r$ and $r+d r$ can also be deduced from the metric:

$$
d V=4 \pi r^{2} \frac{d r}{\sqrt{1-K r^{2}}}
$$


(this formula is for the full sky, for part of the sky we take a fraction)

It's similar to the standard result for spherical polar coordinates, modified for the effect of curvature

## Age of the Universe

- The age of the Universe is the coordinate time $t$ that elapses between $a=0$ and $a=1$
- As an estimate, we can extrapolate today's expansion rate $\frac{d a}{d t}=H_{0}$ :

$$
a(t)=1+H_{0}\left(t-t_{0}\right)
$$

- This leads to an approximate age $\approx$ $\frac{1}{H_{0}}=$ Hubble time $=14 \mathrm{Gyr}$ if $H_{0}=$
 $70 \mathrm{~km} \mathrm{~s}^{-1} \mathrm{Mpc}^{-1}$
- More precisely we would have an integral: $t_{\mathrm{age}}=\int d t=\int_{0}^{1} \frac{d a}{a H(a)}$


## Age of the Universe

- Does an age of $\approx 14$ billion years make sense?

Geological evidence suggests the Earth is $\approx 4.5$ billion years old


Some white dwarf stars are known to be $\approx 12$ billion years old

Radioactive dating shows meteorites are $\approx 4.5$ billion years old


Globular clusters can also be dated as $\approx 12$ billion years old

## Next steps ...

- To go further in our calculations, we'll need to know how the scale factor $a(t)$ depends on time
- This depends on the contents of the Universe, which we'll study next week!




## Key take-aways

- Metrics of a space allow us to relate distance elements to changes in the coordinates of the space
- The Universe is described by combining General Relativity with the assumptions of homogeneity and isotropy
- The result is the Friedmann-Robertson-Walker space-time metric of the expanding Universe, which is written in terms of a constant curvature $K$ and scale factor $a(t)$
- We can measure distances in cosmology using the luminosity distance or angular diameter distance
- The age of the Universe may be deduced as the time elapsed between $a=0$ and $a=1$

