Class 14 : Completing Maxwell's Equations

- Electromagnetism as a synthesis of ideas
- How does Faraday's Law change Maxwell's equations?
- How do the relations between fields and potentials change?
- Displacement current

Recap (1)

• **Maxwell's Equations** for the \vec{E} -field and \vec{B} -field generated by *stationary* (i.e., not varying with time) charge density ρ and current density \vec{J} are:

$$\vec{\nabla} \cdot \vec{E} = \frac{\rho}{\varepsilon_0} \qquad \vec{\nabla} \times \vec{E} = \vec{0}$$
$$\vec{\nabla} \cdot \vec{B} = 0 \qquad \vec{\nabla} \times \vec{B} = \mu_0 \vec{J}$$

Recap (2)

- Extra phenomena are produced by time-varying electromagnetic fields
- Faraday's Law describes how a changing magnetic flux Φ through a circuit induces a voltage in the circuit, $V = -\frac{\partial \Phi}{\partial t}$





 Physics is about the synthesis of ideas : understanding apparently different phenomena as the joint consequences of a deeper reality



 For example, in 1666 Newton realized that the same gravitational force causes both apples to fall downwards from trees on the Earth, and the Earth to orbit the Sun

• Electromagnetism represents another great synthesis of ideas. Prior to 1830, *electricity* and *magnetism* were considered separate phenomena

Electricity







• However, Faraday's experiments demonstrated that a magnet, as well as a battery, can drive a current

Electromagnetism



 Electricity and magnetism are therefore connected by deeper principles



 Charges produce electric fields, and currents produce magnetic fields. But seen in a moving reference frame, a charge becomes a current!



 Special relativity – which describes how physics is viewed in different reference frames – must allow us to transform electrostatics into magnetism

 The other great synthesis of ideas is that electromagnetism can explain optics (light waves)



Faraday's Law

- Faraday's Law $V = -\frac{\partial \Phi}{\partial t}$ requires a modification to Maxwell's equation for electrostatics, $\vec{\nabla} \times \vec{E} = \vec{0}$
- Consider a closed loop *L* bounding a surface *S*



Faraday's Law

- Faraday's Law $V = -\frac{\partial \Phi}{\partial t}$ requires a modification to Maxwell's equation for electrostatics, $\vec{\nabla} \times \vec{E} = \vec{0}$
- Consider a closed loop L bounding a surface S. The potential difference around the loop may be written as $V = \oint \vec{E} \cdot d\vec{l} = \int (\vec{\nabla} \times \vec{E}) \cdot d\vec{A}$, applying Stoke's theorem
- Since $\Phi = \int \vec{B} \cdot d\vec{A}$, Faraday's law can then be written in the form $\int \left(\vec{\nabla} \times \vec{E} + \frac{\partial \vec{B}}{\partial t} \right) \cdot d\vec{A} = 0$
- This holds true for any surface, hence $\vec{\nabla} \times \vec{E} + \frac{\partial \vec{B}}{\partial t} = \vec{0}$

Faraday's Law

- What is the implication for potentials?
- The new relation $\vec{\nabla} \times \vec{E} + \frac{\partial \vec{B}}{\partial t} = \vec{0}$ implies that we cannot derive the electric field as the gradient of an electrostatic potential, $\vec{E} = -\vec{\nabla}V$, for time-varying situations
- However, by substituting in the magnetic vector potential $\vec{B} = \vec{\nabla} \times \vec{A}$, we can derive $\vec{E} = -\vec{\nabla}V \frac{\partial \vec{A}}{\partial t}$
- This is a nice way to describe how electric fields are generated by both *electrostatic charges* $(-\vec{\nabla}V)$ and *changing magnetic fields* $(-\frac{\partial \vec{A}}{\partial t})$

- Maxwell's Equations must also be modified in the presence of time-varying electric fields. The missing ingredient can be illustrated in a couple of ways.
- First, consider taking the divergence of Ampere's Law $\vec{\nabla} \times \vec{B} = \mu_0 \vec{J}$. Since $\vec{\nabla} \cdot (\vec{\nabla} \times \vec{B}) = 0$, we find $\vec{\nabla} \cdot \vec{J} = 0$
- However, charge conservation that $\vec{\nabla} \cdot \vec{J} + \frac{\partial \rho}{\partial t} = 0$
- This contradiction implies that Ampere's law must be wrong for time-varying currents!

- This is fixed by the change $\vec{\nabla} \times \vec{B} = \mu_0 \vec{J} + \mu_0 \varepsilon_0 \frac{\partial \vec{E}}{\partial t}$
- Take the divergence of both sides. We obtain $\mu_0 \vec{\nabla} \cdot \vec{J} + \mu_0 \varepsilon_0 \vec{\nabla} \cdot \frac{\partial \vec{E}}{\partial t} = 0$. Cancelling μ_0 and swapping the order of $\vec{\nabla}$ and $\frac{\partial}{\partial t}$, this becomes $\vec{\nabla} \cdot \vec{J} + \varepsilon_0 \frac{\partial}{\partial t} (\vec{\nabla} \cdot \vec{E}) = 0$
- Now substituting in $\vec{\nabla} \cdot \vec{E} = \frac{\rho}{\varepsilon_0}$, we find $\vec{\nabla} \cdot \vec{J} + \frac{\partial \rho}{\partial t} = 0$
- The new version of the equation is consistent with charge conservation!

- For time-varying \vec{B} -fields : $\vec{\nabla} \times \vec{E} + \frac{\partial \vec{B}}{\partial t} = \vec{0}$
- For time-varying \vec{E} -fields : $\vec{\nabla} \times \vec{B} = \mu_0 \vec{J} + \mu_0 \varepsilon_0 \frac{\partial \vec{E}}{\partial t}$
- A changing \vec{B} -field generates an \vec{E} -field, and a changing \vec{E} -field generates a \vec{B} -field
- The term $\varepsilon_0 \frac{\partial \vec{E}}{\partial t}$ in the 2nd equation is known as the **displacement current**, since it acts like a current density and is equal to $\frac{\partial \vec{D}}{\partial t}$, in terms of the electric displacement $\vec{D} = \varepsilon_0 \vec{E}$

• Next, consider a **capacitor discharging into a circuit**:



• Consider applying Ampere's Law $\oint \vec{B} \cdot d\vec{l} = \mu_0 I_{enclosed}$ to both the *plane surface* and the *bulging surface* in the diagram

 The current enclosed by these two surfaces is different – we again have a problem!



- The situation can be reconciled by including the **displacement current** enclosed by the bulging surface
- The electric field $E = \sigma/\varepsilon_0$ in terms of the charge density σ
- If A is the plate area, then $Q = \sigma A$ hence $I = \frac{dQ}{dt} = \varepsilon_0 \frac{dE}{dt} A$
- This is equal to the displacement current enclosed by the surface!

Maxwell's complete equations

- Maxwell's Equations are now complete!
- In a vacuum:

$$\vec{\nabla} \cdot \vec{E} = \frac{\rho}{\varepsilon_0} \qquad \vec{\nabla} \times \vec{E} + \frac{\partial \vec{B}}{\partial t} = \vec{0}$$
$$\vec{\nabla} \cdot \vec{B} = 0 \qquad \vec{\nabla} \times \vec{B} - \mu_0 \varepsilon_0 \frac{\partial \vec{E}}{\partial t} = \mu_0 \vec{J}$$

Maxwell's complete equations

• More generally, in materials with permittivity ε_r and permeability μ_r we can use the electric displacement $\vec{D} = \varepsilon_r \varepsilon_0 \vec{E}$ and magnetic intensity $\vec{H} = \vec{B}/\mu_r \mu_0$:

$$\vec{\nabla} \cdot \vec{D} = \rho \qquad \qquad \vec{\nabla} \times \vec{E} + \frac{\partial \vec{B}}{\partial t} = \vec{0}$$
$$\vec{\nabla} \cdot \vec{B} = 0 \qquad \qquad \vec{\nabla} \times \vec{H} - \frac{\partial \vec{D}}{\partial t} = \vec{J}$$

Ampere's Law $\oint \vec{B} \cdot d\vec{l} = \mu_0 I_{enclosed}$ holds true ...

- A. Always
- B. When the magnetic field \vec{B} is not changing
- C. When the electric field \vec{E} is not changing
- D. When both \vec{E} and \vec{B} are not changing



Gauss's Law $\int \vec{E} \cdot d\vec{A} = \frac{Q_{enclosed}}{\varepsilon_0}$ holds true ...

- A. Always
- B. When the magnetic field \vec{B} is not changing
- C. When the electric field \vec{E} is not changing
- D. When both \vec{E} and \vec{B} are not changing



An electrostatic potential $V = -\int_{\infty}^{P} \vec{E} \cdot d\vec{l}$ can be defined ...

- A. Always
- B. When the magnetic field \vec{B} is not changing
- C. When the electric field \vec{E} is not changing
- D. When both \vec{E} and \vec{B} are not changing





- A. U
- B. = $\mu_0 I$
- C. = $-\mu_0 I$
- D. Impossible to say without knowing surface



Summary

- An electric field is generated by a changing magnetic field
 (electromagnetic induction)
- A magnetic field is also generated by a changing electric field, as described by the **displacement current** such that $\vec{\nabla} \times \vec{B} = \mu_0 \vec{J} + \mu_0 \varepsilon_0 \frac{\partial \vec{E}}{\partial t}$

$$\nabla \cdot \mathbf{E} = \frac{\rho}{\varepsilon_0}$$
$$\nabla \cdot \mathbf{B} = 0$$
$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$
$$\nabla \times \mathbf{B} = \mu_0 \mathbf{J} + \mu_0 \varepsilon_0 \frac{\partial \mathbf{E}}{\partial t}$$